



# *Exchange Energy Formulations for Micromagnetics*

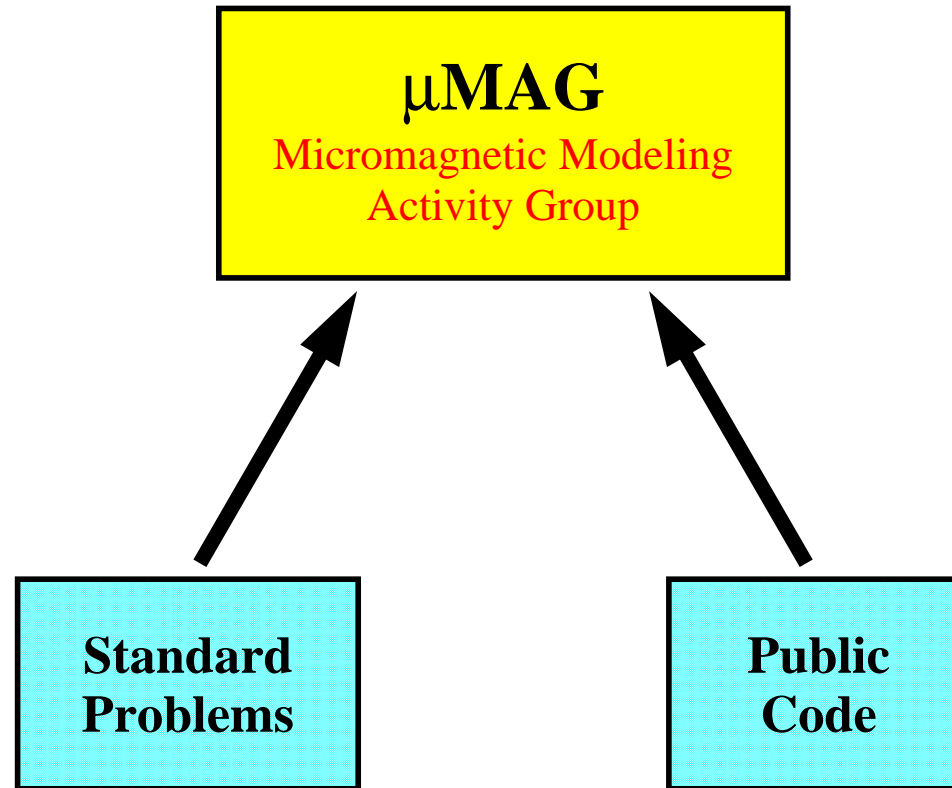
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# Outline

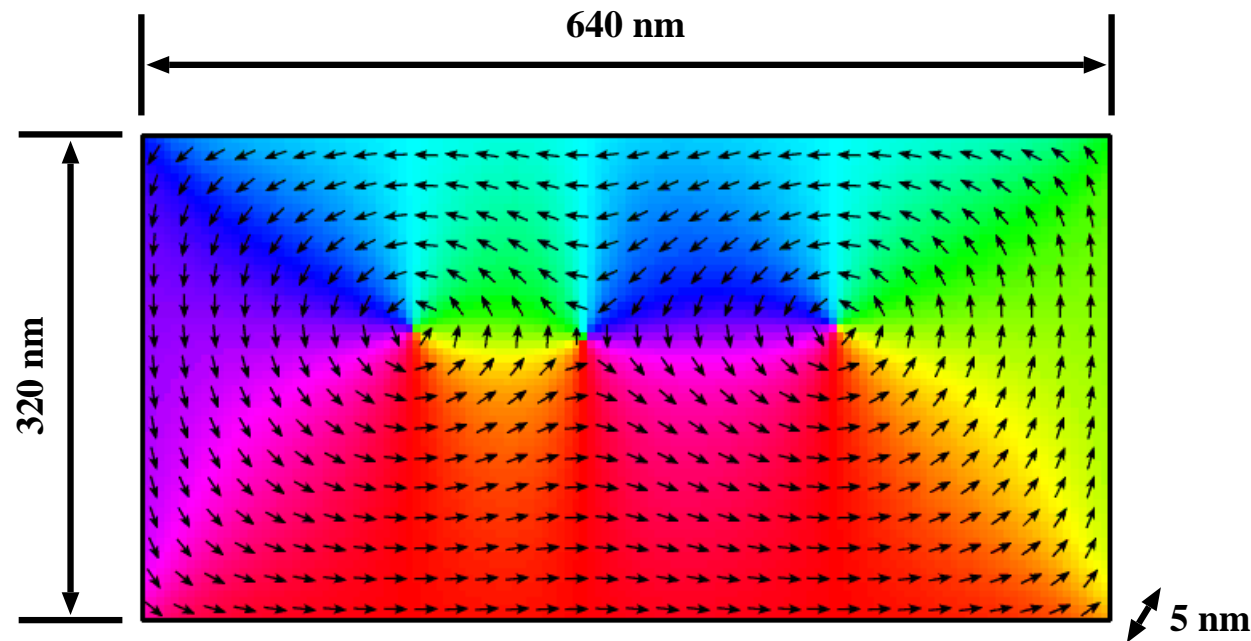
- Background
  - Quasi-static vs. Dynamic
- Exchange energy formulations
  - 6-ngbr, 12-ngbr, 26-ngbr
  - Numerical integration
  - Integrand representation
  - Boundary conditions



Center for Theoretical and Computational Materials Science

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# Micromagnetics



The study, modeling and simulation of magnetic materials and their behavior at the nanometer scale.

# Brown's equations

## Energies:

$$E_{\text{exchange}} = \int_V \frac{A}{M_s^2} (|\nabla M_x|^2 + |\nabla M_y|^2 + |\nabla M_z|^2) d^3r$$

$$E_{\text{anisotropy}} = \int_V \frac{K_1}{M_s^2} (\mathbf{M} \cdot \mathbf{u})^2 d^3r$$

$$E_{\text{demag}} = \frac{\mu_0}{8\pi} \int_V \mathbf{M}(\mathbf{r}) \cdot \left[ \int_V \nabla \cdot \mathbf{M}(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|^3} d^3r' - \int_S \hat{\mathbf{n}} \cdot \mathbf{M}(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|^3} d^2r' \right] d^3r$$

$$E_{\text{Zeeman}} = -\mu_0 \int_V \mathbf{M} \cdot \mathbf{H}_{\text{ext}} d^3r$$

# Constraints

$\mathbf{M}$  is smooth

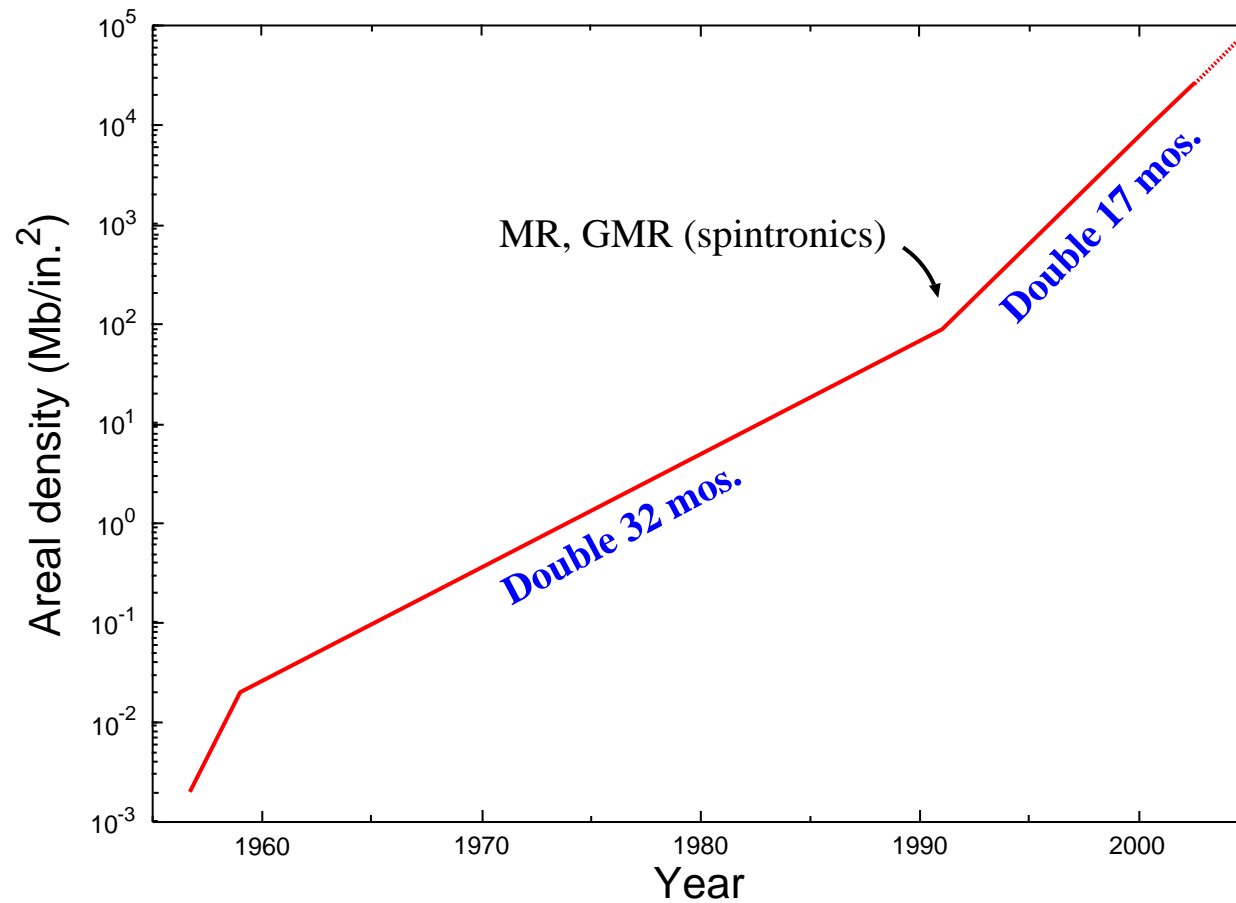
and

$$\|\mathbf{M}(t)\| = \|\mathbf{M}\| = M_s$$

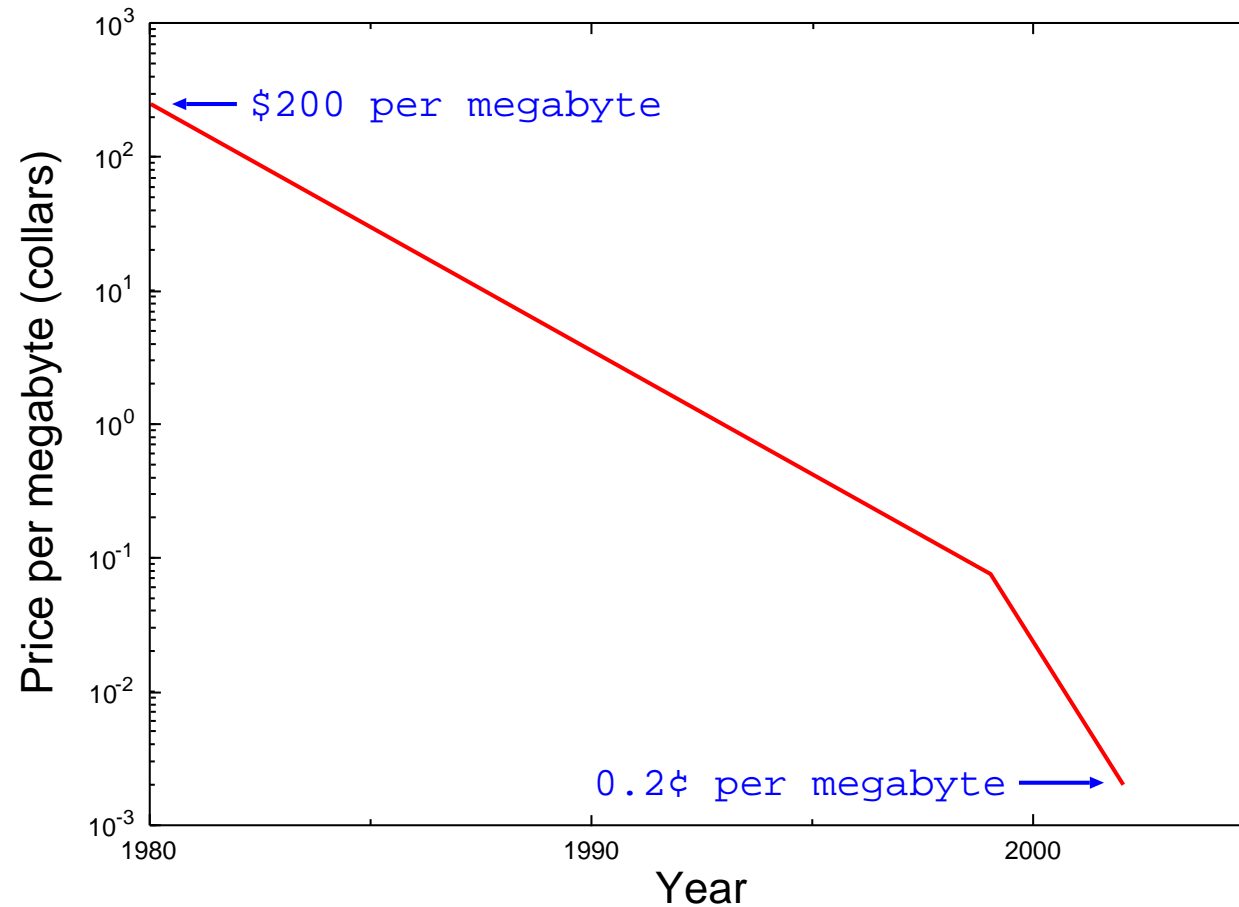
or equivalently

$$\|\mathbf{m}(t)\| = \|\mathbf{m}\| = \mathbf{M}/M_s = 1.$$

# Magnetic disk storage

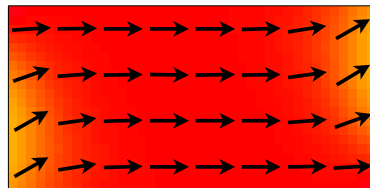
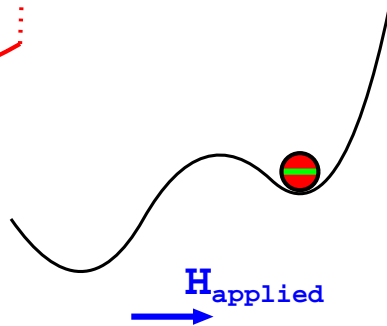
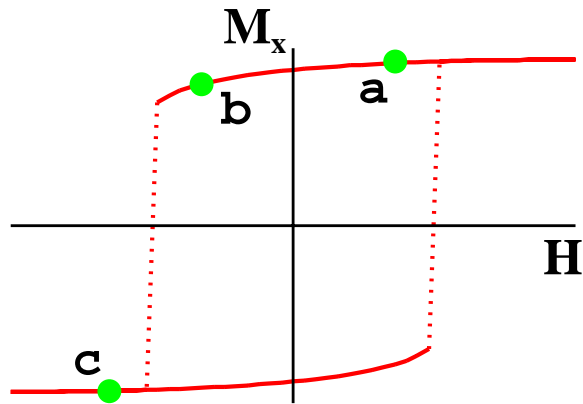


# Magnetic disk storage

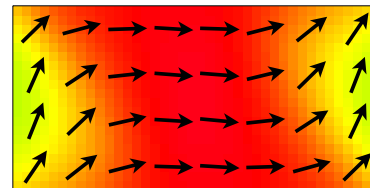
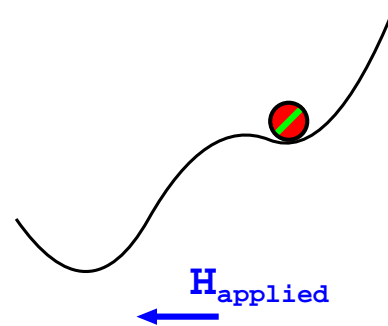




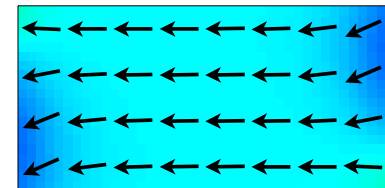
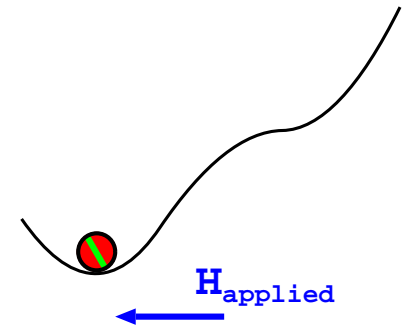
# Quasi-static micromagnetics



a



b



c

# Magnetization dynamics

## Landau-Lifshitz-Gilbert:

$$\frac{d\mathbf{M}}{dt} = \frac{-\omega}{1 + \alpha^2} \mathbf{M} \times \mathbf{H}_{\text{eff}} - \frac{\alpha \omega}{(1 + \alpha^2) M_s} \mathbf{M} \times (\mathbf{M} \times \mathbf{H}_{\text{eff}})$$

where

$$\mathbf{H}_{\text{eff}} = -\frac{1}{\mu_0} \frac{\partial E}{\partial \mathbf{M}}$$

$\omega$  = gyromagnetic ratio

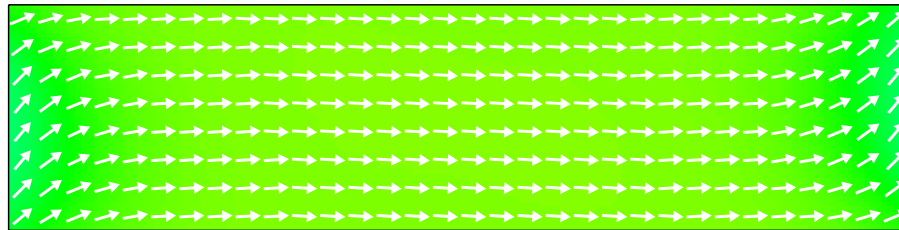
$\alpha$  = damping coefficient

# Magnetization dynamics

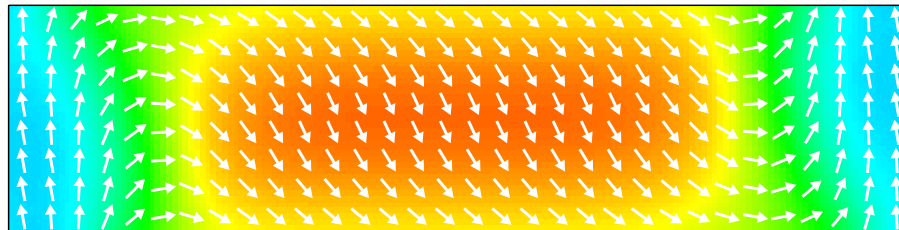
Time

$\mu_0 H = 36 \text{ mT}$   $x$

0 ps



100 ps



150 ps

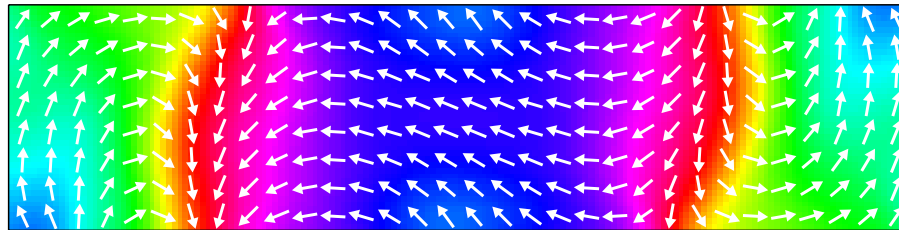


# Magnetization dynamics

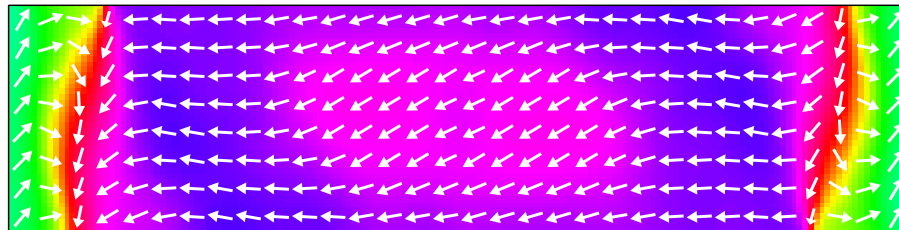
**Time**

$\mu_0 H = 36 \text{ mT}$   $x$

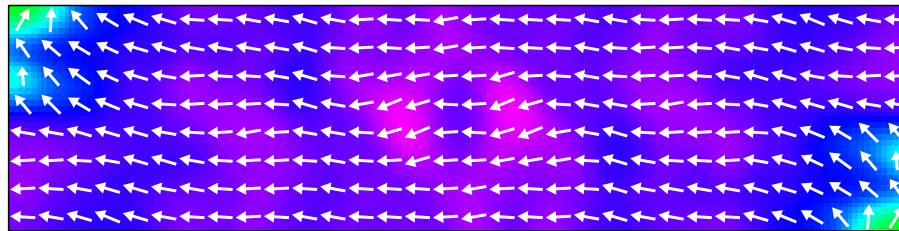
350 ps



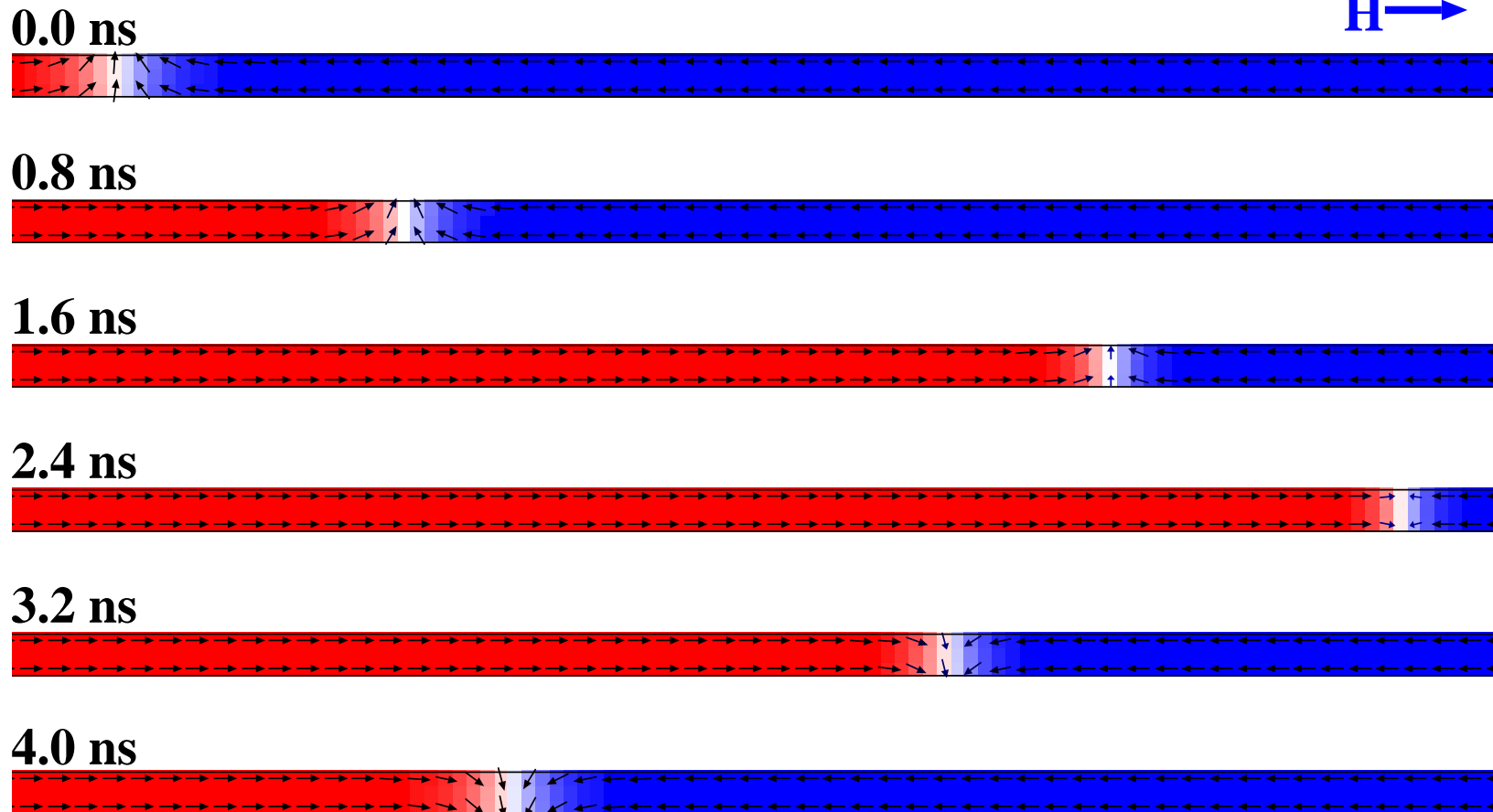
450 ps



750 ps



# Magnetization yoyo





# Variational derivatives

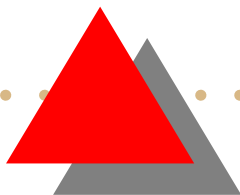
Let  $(\text{support } \Delta \mathbf{M}) \subset B(x_k, \epsilon)$ .

Then

$$\left. \frac{\delta E}{\delta \mathbf{M}} \right|_{x_k} = \lim \frac{E(\mathbf{M} + \Delta \mathbf{M}) - E(\mathbf{M})}{\|\Delta \mathbf{M}\|_1}$$

as

$$\epsilon \rightarrow 0, \quad \|\Delta \mathbf{M}\|_\infty \rightarrow 0.$$





# Variational derivatives

In particular, if

$$\mathbf{M}(x) = \sum \mathbf{M}_i \phi_i(x),$$

then

$$\left. \frac{\delta E}{\delta \mathbf{M}} \right|_{x_k} \approx \frac{\partial E}{\partial \mathbf{M}_k} \cdot \frac{1}{\|\phi_k\|_1}$$

# Brown's equations

## Energies:

$$E_{\text{exchange}} = \int_V \frac{A}{M_s^2} (|\nabla M_x|^2 + |\nabla M_y|^2 + |\nabla M_z|^2) d^3r$$

$$E_{\text{anisotropy}} = \int_V \frac{K_1}{M_s^2} (\mathbf{M} \cdot \mathbf{u})^2 d^3r$$

$$E_{\text{demag}} = \frac{\mu_0}{8\pi} \int_V \mathbf{M}(\mathbf{r}) \cdot \left[ \int_V \nabla \cdot \mathbf{M}(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|^3} d^3r' - \int_S \hat{\mathbf{n}} \cdot \mathbf{M}(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|^3} d^2r' \right] d^3r$$

$$E_{\text{Zeeman}} = -\mu_0 \int_V \mathbf{M} \cdot \mathbf{H}_{\text{ext}} d^3r$$



# Discrete approximation

$$\begin{aligned} E_{\text{exchange}} &= \int_V A (|\nabla m_x|^2 + |\nabla m_y|^2 + |\nabla m_z|^2) d^3r \\ &= \Phi[\mathbf{m}(\mathbf{x}_1), \mathbf{m}(\mathbf{x}_2), \dots, \mathbf{m}(\mathbf{x}_n)] + O(h^k) \end{aligned}$$

where

$h$  is cell size

$k$  is approximation order

# *Discrete approximation*

$$E_{\text{exchange}} = \int_V A (|\nabla m_x|^2 + |\nabla m_y|^2 + |\nabla m_z|^2) d^3r$$

- Numerical integration
- Integrand representation
- Boundary conditions

# Numerical integration

$$\int_a^b f \approx h \sum w_k f_k$$

Closed intervals,  $x_k = a + kh$ ,

$$O(h^2) \text{ error: } (w_k) = \left[ \frac{1}{2} \ 1 \ 1 \ \dots \ 1 \ \frac{1}{2} \right]$$

$$O(h^4) \text{ error: } (w_k) = \frac{1}{3} \left[ 1 \ 4 \ 2 \ 4 \ \dots \ 2 \ 4 \ 1 \right]$$

$$O(h^4) \text{ error: } (w_k) = \left[ \frac{3}{8} \ \frac{7}{6} \ \frac{23}{24} \ 1 \ 1 \ \dots \ 1 \ \frac{23}{24} \ \frac{7}{6} \ \frac{3}{8} \right]$$

# Numerical integration

$$\int_a^b f \approx h \sum w_k f_k$$

Open intervals,  $x_k = a + (k - 1/2)h$ ,

$$O(h^2) \text{ error: } (w_k) = [1 \ 1 \ 1 \ \dots \ 1]$$

$$O(h^4) \text{ error: } (w_k) = \left[ \frac{13}{12} \ \frac{7}{8} \ \frac{25}{24} \ 1 \ 1 \ \dots \ 1 \ \frac{25}{24} \ \frac{7}{8} \ \frac{13}{12} \right]$$

# Integrand representation

$$\begin{aligned} E_{\text{exchange}} &= \int_V A (|\nabla m_x|^2 + |\nabla m_y|^2 + |\nabla m_z|^2) d^3r \\ &= - \int_V A \mathbf{m} \cdot \left( \frac{\partial^2 \mathbf{m}}{\partial x^2} + \frac{\partial^2 \mathbf{m}}{\partial y^2} + \frac{\partial^2 \mathbf{m}}{\partial z^2} \right) d^3r \end{aligned}$$

Since

$$|\nabla f|^2 = \nabla \cdot (f \nabla f) - f \nabla^2 f$$

and

$$\|\mathbf{m}\| = 1.$$

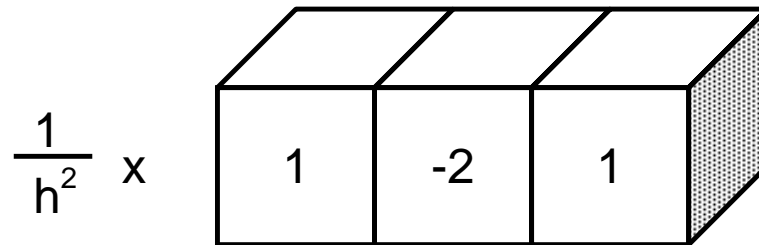
# Discretized energy

$$- \iiint A \mathbf{m} \cdot \frac{\partial^2 \mathbf{m}}{\partial x^2} d^3 r$$

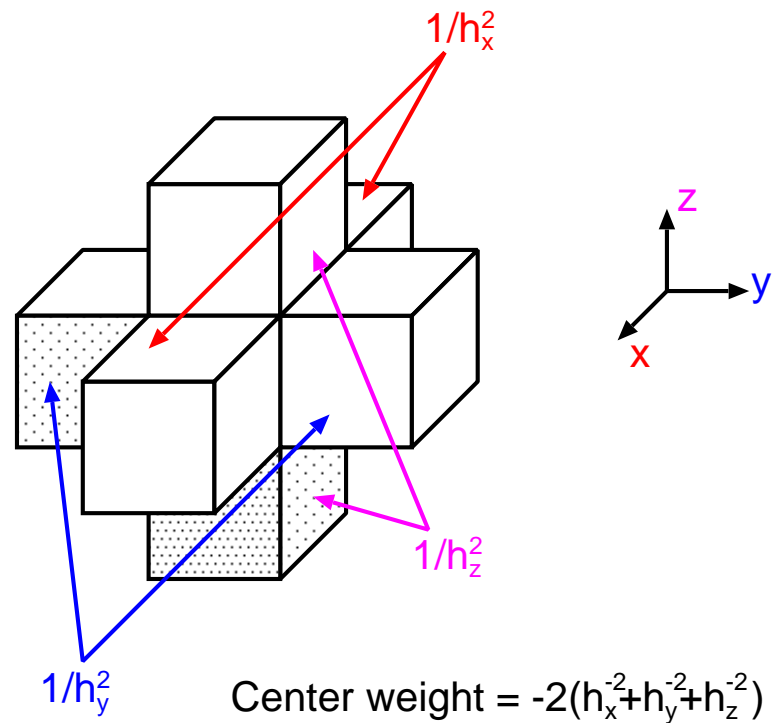
$$\approx -h_x h_y h_z \sum_{kjii'} w_k^z w_j^y w_i^x A_{ijk} d_{ii'} \mathbf{m}_{ijk} \cdot \mathbf{m}_{i'jk}$$

# 3-pt stencil

$$\frac{\partial^2 \mathbf{m}(x)}{\partial x^2} = \frac{1}{h^2} [\mathbf{m}(x - h) - 2\mathbf{m}(x) + \mathbf{m}(x + h)] + O(h^2)$$



# 3-pt stencil

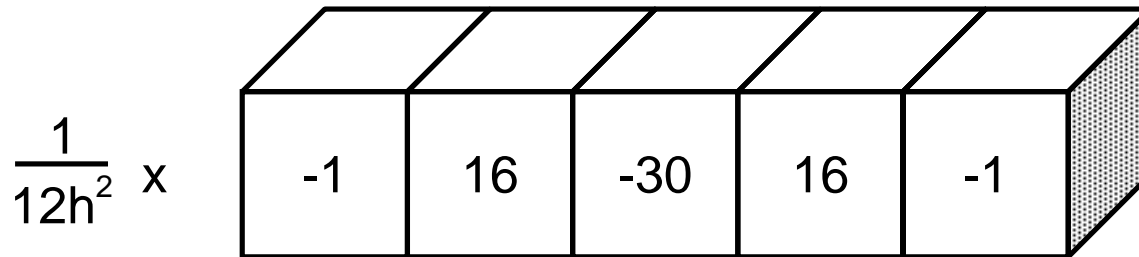


“6-neighbor exchange”

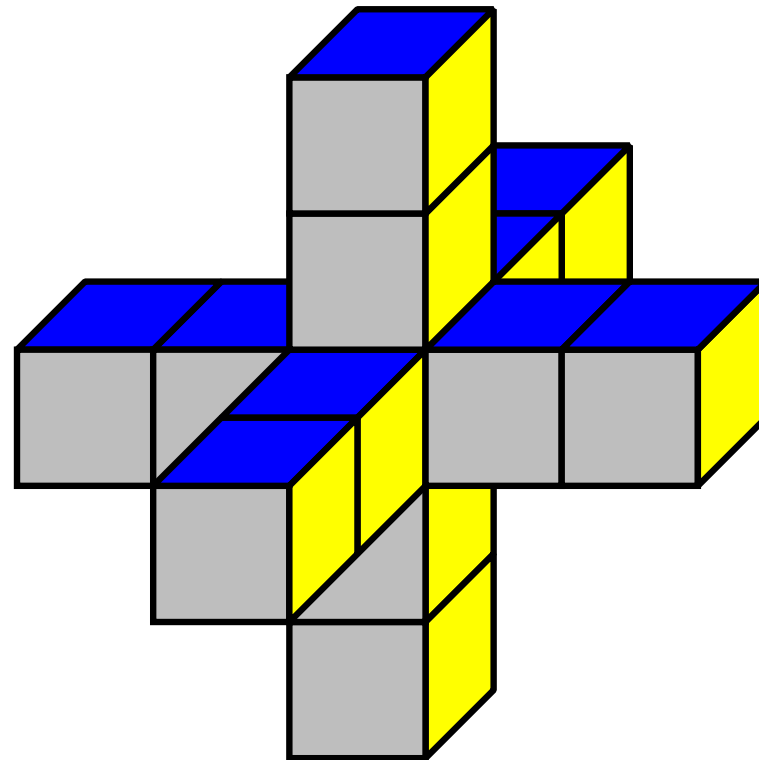


# 5-pt stencil

$$\frac{\partial^2 \mathbf{m}(x)}{\partial x^2} = \frac{1}{12h^2} [-\mathbf{m}(x - 2h) + 16\mathbf{m}(x - h) - 30\mathbf{m}(x) + 16\mathbf{m}(x + h) - \mathbf{m}(x + 2h)] + O(h^4)$$

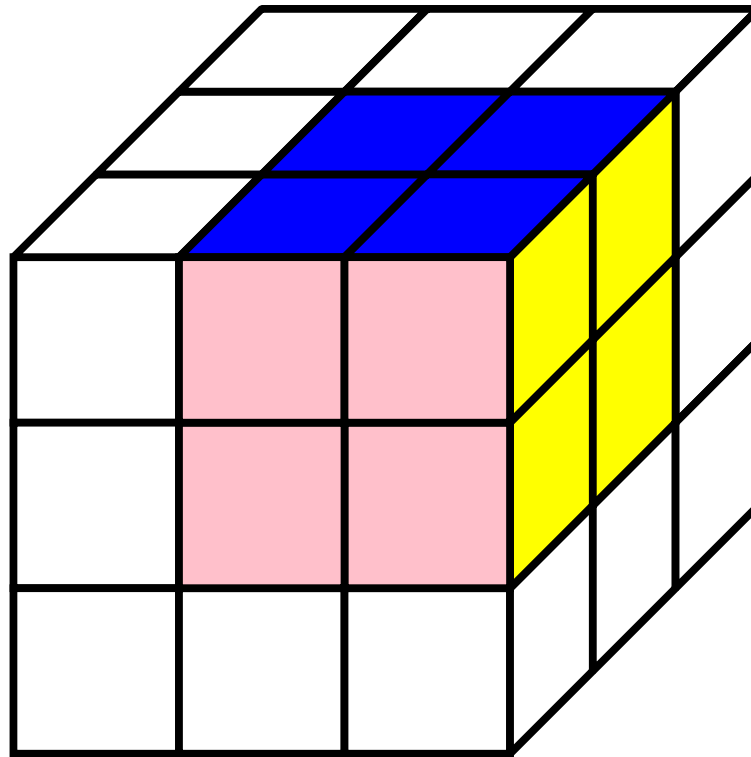


# *5-pt stencil*



“12-neighbor exchange”

# *Trilinear interpolation*



“26-neighbor exchange”

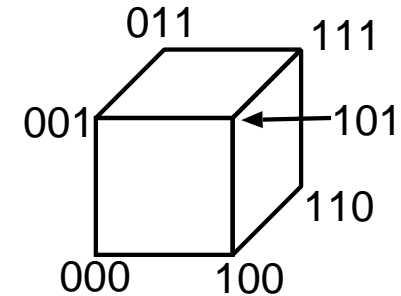
# Trilinear interpolation

Given  $\mathbf{m}_{000}$ ,  $\mathbf{m}_{100}$ ,  $\dots$ , solve for

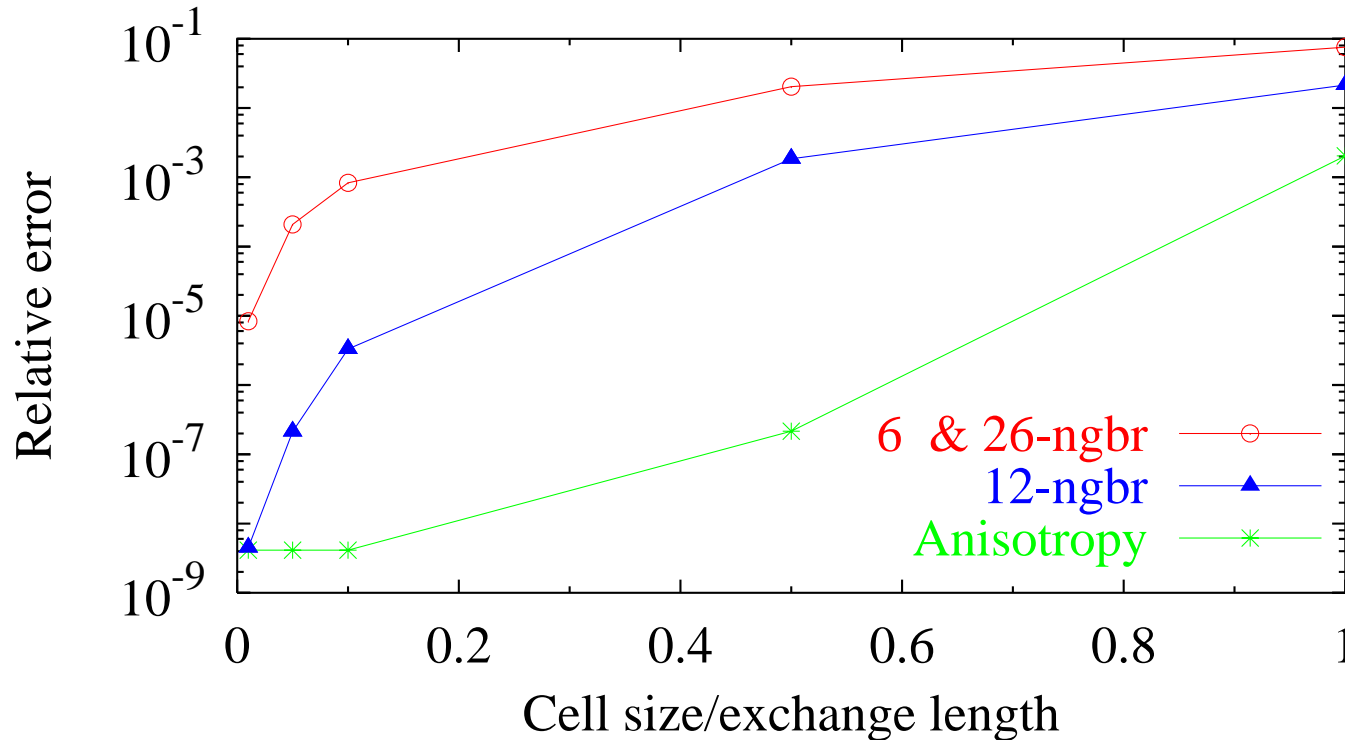
$$\begin{aligned} \mathbf{m}(x) = & \mathbf{a}_0 + \mathbf{a}_{100}x + \mathbf{a}_{010}y + \mathbf{a}_{001}z \\ & + \mathbf{a}_{110}xy + \mathbf{a}_{101}xz + \mathbf{a}_{011}yz + \mathbf{a}_{111}xyz. \end{aligned}$$

Then use

$$E_{\text{exchange}} = \int_V A (|\nabla m_x|^2 + |\nabla m_y|^2 + |\nabla m_z|^2) d^3r$$



# Analytic 1D domain wall



Relative energy error vs. discretization cell size

# *Exchange lengths*

Magnetostatic-exchange length  $= \sqrt{\frac{2A}{\mu_0 M_s^2}}$

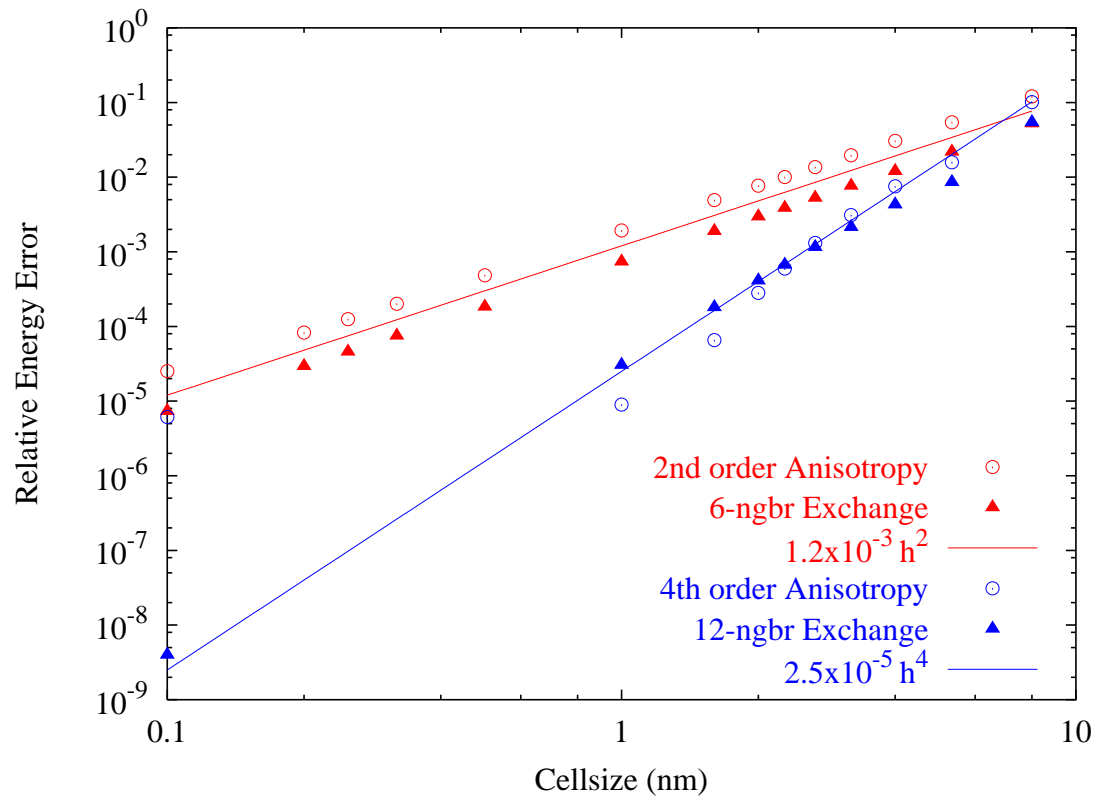
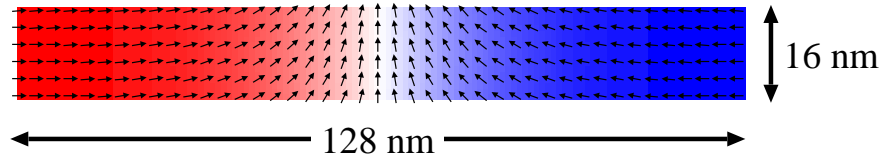
Magnetocrystalline-exchange length  $= \sqrt{\frac{A}{|K|}}$

# Convergence of equilibrium

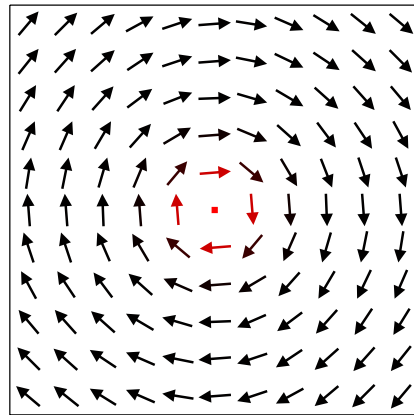
$$A=13 \times 10^{-12} \text{ J/m}$$

$$K_u=(4.6 \times 10^5)r^2/(1+r^2) \text{ J/m}^3$$

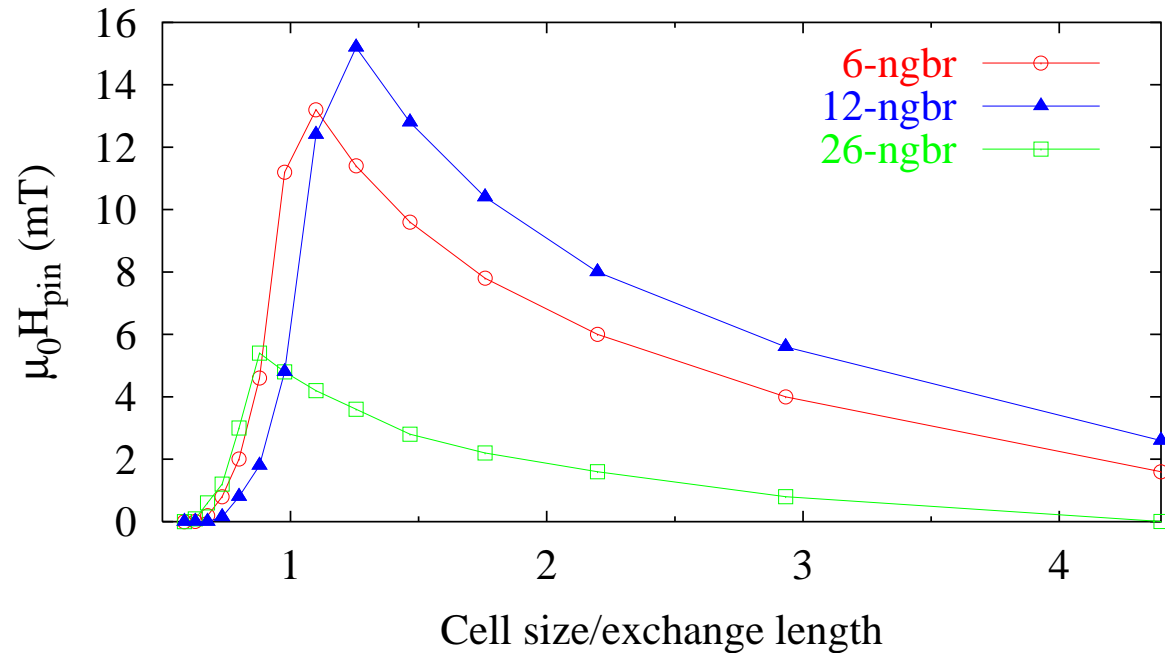
10 nm thick



# Vortex mobility



$H_{\text{applied}}$   $\rightarrow$  Vortex motion



(Compare to Donahue & McMichael, Physica B, **233**, 272 (1997).)



# Discretized energy

$$- \iiint A \mathbf{m} \cdot \frac{\partial^2 \mathbf{m}}{\partial x^2} d^3 r$$

$$\approx -h_x h_y h_z \sum_{kjii'} w_k^z w_j^y w_i^x A_{ijk} d_{ii'} \mathbf{m}_{ijk} \cdot \mathbf{m}_{i'jk}$$

# Boundary?

$$\frac{1}{12h^2} \times \begin{bmatrix} \boxed{?} & \boxed{?} \\ & \boxed{?} \end{bmatrix} \begin{bmatrix} -30 & 16 & -1 & & & & \\ 16 & -30 & 16 & -1 & & & \\ -1 & 16 & -30 & 16 & -1 & & \\ & -1 & 16 & -30 & 16 & -1 & \\ & & & & \ddots & & \end{bmatrix}$$

# Variational calculus

Let

$$E[m] = \int_a^b f(x, m, m') dx$$

Then

$$\begin{aligned} E[m + h] - E[m] &= \int_a^b \left( f_m - \frac{d}{dx} f_{m'} \right) h dx \\ &+ h(b) f_{m'}(b, m(b), m'(b)) - h(a) f_{m'}(a, m(a), m'(a)) \\ &+ O(h^2 + h'^2). \end{aligned}$$



# *Euler-Lagrange eqn*

If  $m$  is extremal, then

$$f_m - \frac{d}{dx} f_{m'} = 0 \quad (\text{Euler-Lagrange})$$



# Boundary conditions

Since

$$h(a) f_{m'}(a, m(a), m'(a)) = 0,$$

if  $m(a)$  is free, then

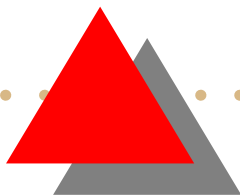
$$f_{m'}(a, m(a), m'(a)) = 0.$$

But

$$f(x, m, m') = Am'^2 + g(x, m)$$

and

$$f_{m'} = 2Am' \Rightarrow m'(a) = 0.$$



# Boundary?

$$\frac{1}{h^2} \times \boxed{?} \begin{bmatrix} -2 & 1 & & & \\ 1 & -2 & 1 & & \\ & 1 & -2 & 1 & \\ & & 1 & -2 & 1 \\ & & & \ddots & \ddots \end{bmatrix}$$

# Boundary conditions, 6-ngbr

Neumann boundary:

$$\left. \frac{\partial^2 \mathbf{m}}{\partial x^2} \right|_{x_1} = \frac{\mathbf{m}_2 - \mathbf{m}_1}{h^2} - \frac{1}{h} \left. \frac{\partial \mathbf{m}}{\partial x} \right|_a + O(h).$$

Dirichlet boundary:

$$\left. \frac{\partial^2 \mathbf{m}}{\partial x^2} \right|_{x_1} = \frac{4\mathbf{m}_2 - 12\mathbf{m}_1}{3h^2} + \frac{8}{3h^2} \mathbf{m}(a) + O(h).$$

# 6-ngbr, Neumann

$$(d_{ii'}) = \frac{1}{h^2} \begin{bmatrix} -1 & 1 & & & & \\ 1 & -2 & 1 & & & \\ & 1 & -2 & 1 & & \\ & & 1 & -2 & 1 & \\ & & & 1 & -2 & 1 \\ & & & & \ddots & \ddots \end{bmatrix}$$

+ boundary derivative field (if any).



# 6-ngbr, Dirichlet

$$(d_{ii'}) = \frac{1}{h^2} \begin{bmatrix} -4 & 4/3 & & & & \\ & 1 & -2 & 1 & & \\ & & 1 & -2 & 1 & \\ & & & 1 & -2 & 1 \\ & & & & \ddots & \\ & & & & & \ddots & \end{bmatrix}$$

+ boundary value field.

# Discretized representation

$$\begin{aligned} & - \iiint A \mathbf{m} \cdot \frac{\partial^2 \mathbf{m}}{\partial x^2} d^3 r \\ & \approx -h_x h_y h_z \sum_{jk} w_k^z w_j^y \sum_{ii'} A_{ijk} w_i^x d_{ii'} \mathbf{m}_{ijk} \cdot \mathbf{m}_{i'jk} \\ & \stackrel{\text{def}}{=} \Phi \end{aligned}$$

# Discretized representation

$$\frac{\partial \Phi}{\partial \mathbf{m}_{ijk}} = -2h_x h_y h_z \sum_{jk} w_k^z w_j^y \sum_{ii'} c_{ii'jk} \mathbf{m}_{i'jk}$$

where

$$c_{ii'jk} = (A_{ijk} w_i^x d_{ii'} + A_{i'jk} w_{i'}^x d_{i'i}) / 2$$

or

$$c_{ii'} = A (w_i^x d_{ii'} + w_{i'}^x d_{i'i}) / 2$$

if  $A$  is constant.



# 6-ngbr, Dirichlet

Clean up representation:

- Include  $w_i^x$  terms
- Symmetrize
- Adjust diagonal so row sums = 0

# 6-ngbr, Dirichlet

$$C = \frac{A}{h^2} \begin{bmatrix} -7/6 & 7/6 & & & & \\ 7/6 & -13/6 & 1 & & & \\ & 1 & -2 & 1 & & \\ & & 1 & -2 & 1 & \\ & & & \ddots & \ddots & \\ & & & & & \ddots & \ddots \end{bmatrix}$$

+ boundary value field.

# 12-ngbr, Neumann boundary

$$\frac{\partial^2 \mathbf{m}}{\partial x^2} \Big|_{x_1} = \frac{-59\mathbf{m}_1 + 64\mathbf{m}_2 - 5\mathbf{m}_3}{38h^2} - \frac{16}{19h} \frac{\partial \mathbf{m}}{\partial x} \Big|_a - \frac{11}{19h} \frac{\partial \mathbf{m}}{\partial x} \Big|_{x_1} + O(h^3)$$

Norm constraint  $\Rightarrow \mathbf{m}_1 \cdot \frac{\partial \mathbf{m}}{\partial x} \Big|_{x_1} = 0.$

# 12-ngbr, Neumann boundary

$$\frac{\partial^2 \mathbf{m}}{\partial x^2} \Big|_{x_2} = \frac{335\mathbf{m}_1 - 669\mathbf{m}_2 + 357\mathbf{m}_3 - 23\mathbf{m}_4}{264h^2} + \frac{1}{11h} \frac{\partial \mathbf{m}}{\partial x} \Big|_a + O(h^3)$$



# *12-ngbr, Neumann boundary*

$$C = \frac{A}{12h^2} \begin{bmatrix} -16.2 & 17.6 & -1.4 & & & & & & & & & & & \\ & 17.6 & -32.1 & 15.4 & -1.0 & & & & & & & & & & \\ & -1.4 & 15.4 & -29.4 & 16.3 & -1 & & & & & & & & & \\ & & -1.0 & 16.3 & -30.4 & 16 & -1 & & & & & & & & \\ & & & -1 & 16 & -30 & 16 & -1 & & & & & & & \\ & & & & & & & & \ddots & & & & & & \\ & & & & & & & & & & & & & & & \\ & & & & & & & & & & & & & & & \\ & & & & & & & & & & & & & & & \\ & & & & & & & & & & & & & & & \\ & & & & & & & & & & & & & & & \\ & & & & & & & & & & & & & & & \\ & & & & & & & & & & & & & & & \end{bmatrix} + O(h^4).$$



# 12-ngbr, Dirichlet boundary

$$\frac{\partial^2 \mathbf{m}}{\partial x^2} \Big|_{x_1} = \frac{-165\mathbf{m}_1 + 40\mathbf{m}_2 - 3\mathbf{m}_3}{30h^2} + \frac{64}{15h} \mathbf{m}(a) + \frac{1}{h} \frac{\partial \mathbf{m}}{\partial x} \Big|_{x_1} + O(h^3)$$

Norm constraint  $\Rightarrow \mathbf{m}_1 \cdot \frac{\partial \mathbf{m}}{\partial x} \Big|_{x_1} = 0.$

# 12-ngbr, Dirichlet boundary

$$\left. \frac{\partial^2 \mathbf{m}}{\partial x^2} \right|_{x_2} = \frac{4\mathbf{m}_1 - 15\mathbf{m}_2 + 12\mathbf{m}_3 - 4\mathbf{m}_4}{6h^2}$$

$$-\frac{1}{h} \left. \frac{\partial \mathbf{m}}{\partial x} \right|_{x_2} + O(h^3)$$

Norm constraint  $\Rightarrow \mathbf{m}_2 \cdot \left. \frac{\partial \mathbf{m}}{\partial x} \right|_{x_2} = 0.$

# *12-ngbr, Dirichlet boundary*

$$C = \frac{A}{12h^2} \begin{bmatrix} -11.0 & 12.2 & -1.2 & & & & & & \\ 12.2 & -29.6 & 18.8 & -1.4 & & & & & \\ -1.2 & 18.8 & -33.0 & 16.3 & -1 & & & & \\ & -1.4 & 16.3 & -30.0 & 16 & -1 & & & \\ & & -1 & 16 & -30 & 16 & -1 & & \\ & & & & & & & \ddots & \\ & & & & & & & & & & \end{bmatrix} + O(h^4).$$



# *Eigenvalue analysis*

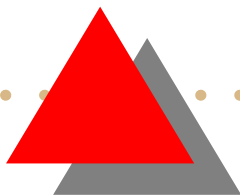
For 6-ngbr methods:

Eigenvalues of  $-C \subset [0, 4)$

For 12-ngbr methods:

Eigenvalues of  $-C \subset [0, 5\frac{1}{3})$

$\Rightarrow$  Good iterative convergence!



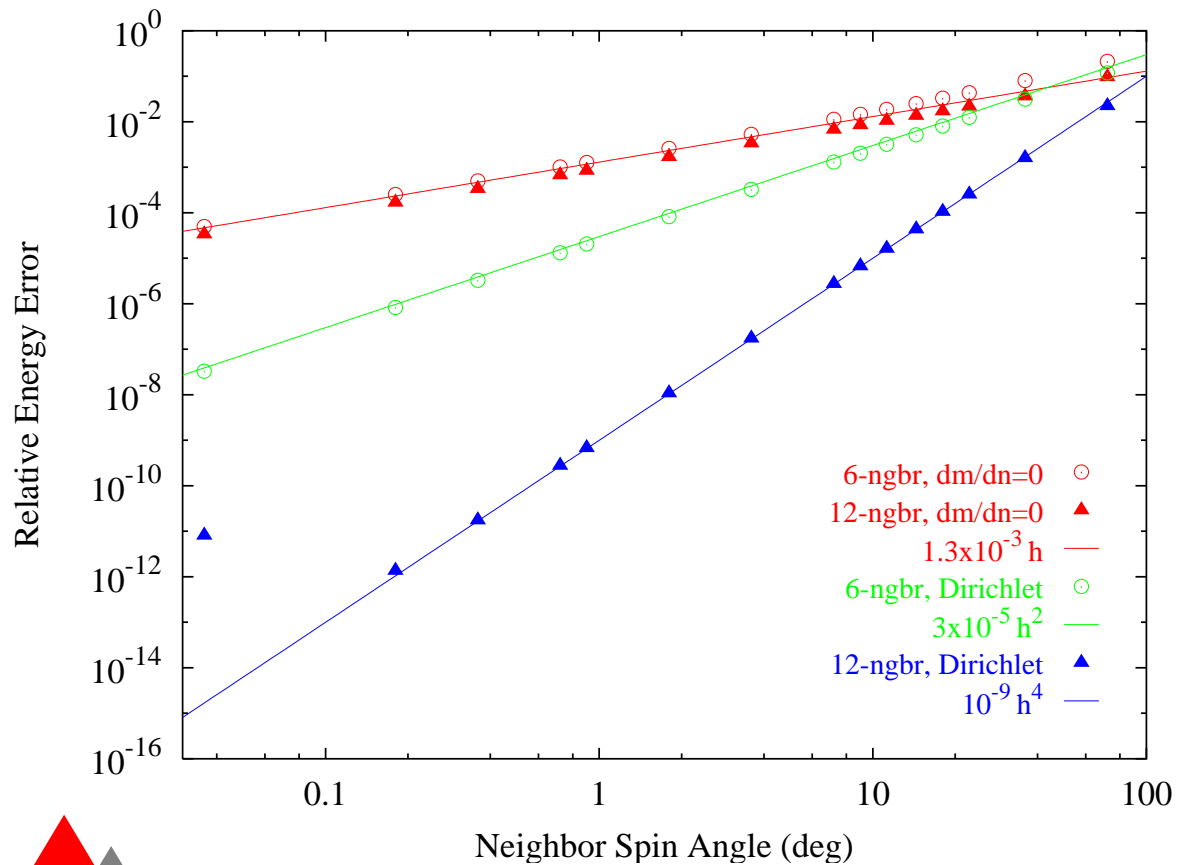
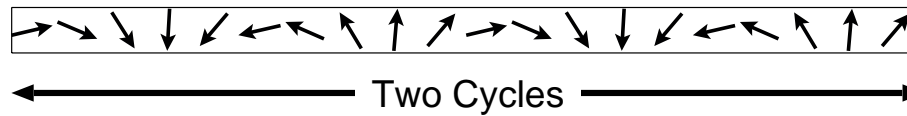
# 12-ngbr exchange

Assume  $\partial \mathbf{m} / \partial \hat{\mathbf{n}} = 0$ :

$$\frac{1}{12h^2} \times \begin{bmatrix} * & * & * & * & & & \\ * & * & * & * & & & \\ -1 & 16 & -30 & 16 & -1 & & \\ & -1 & 16 & -30 & 16 & -1 & \\ & & & \ddots & & & \\ & & & & & & \end{bmatrix}$$

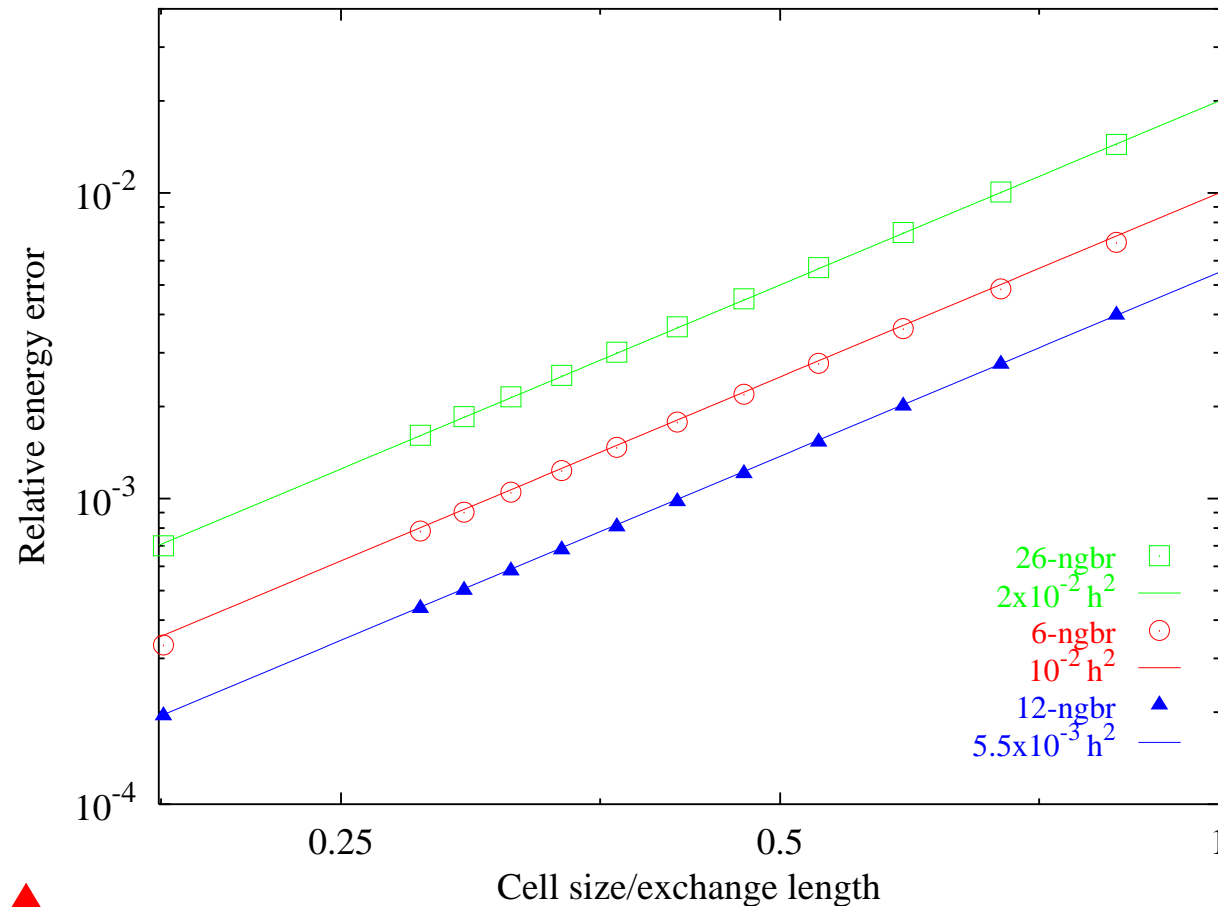
$\Rightarrow$  Bad eigenvalues!

# Magnetization spiral



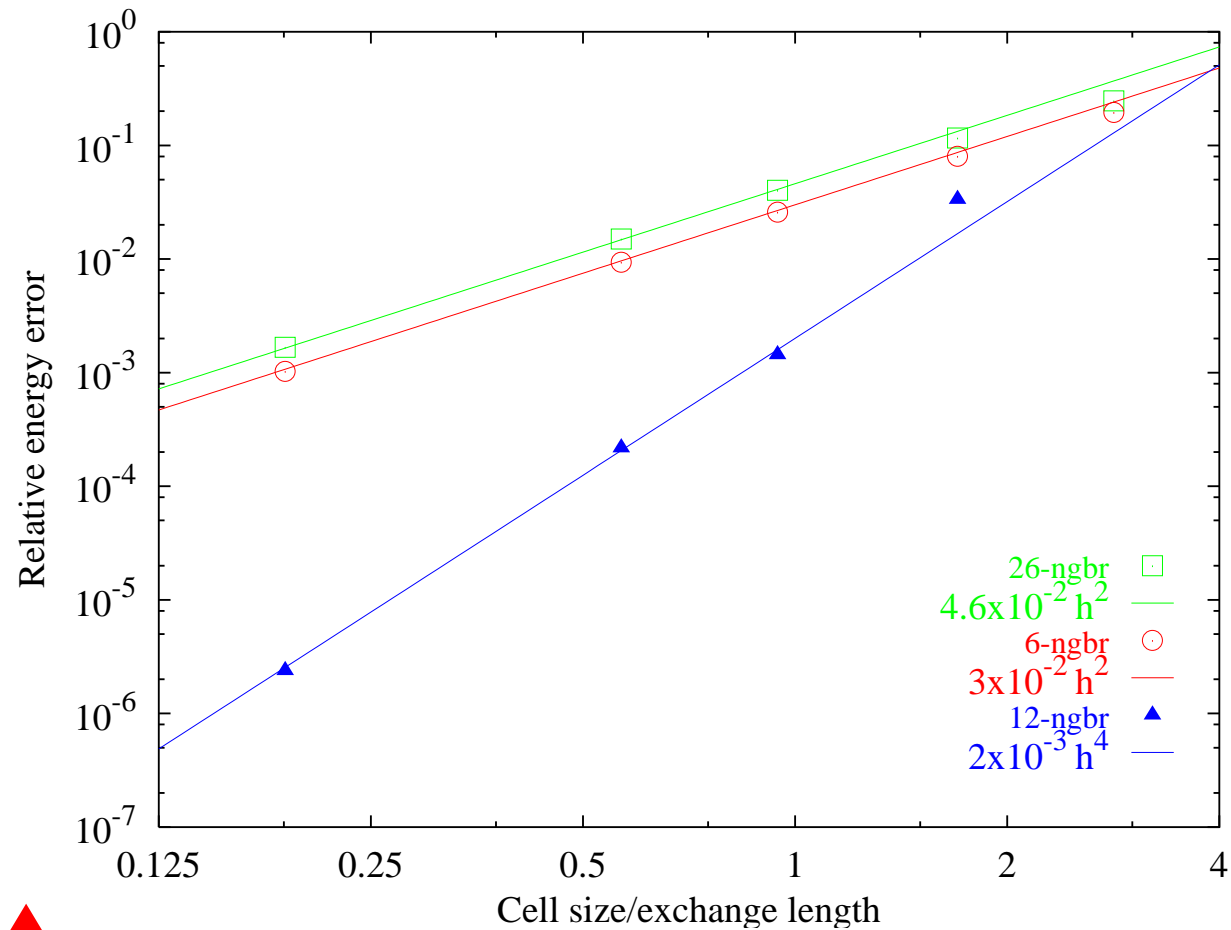
# *muMAG Standard Problem 3*

Equilibria convergence:



# *muMAG Standard Problem 3*

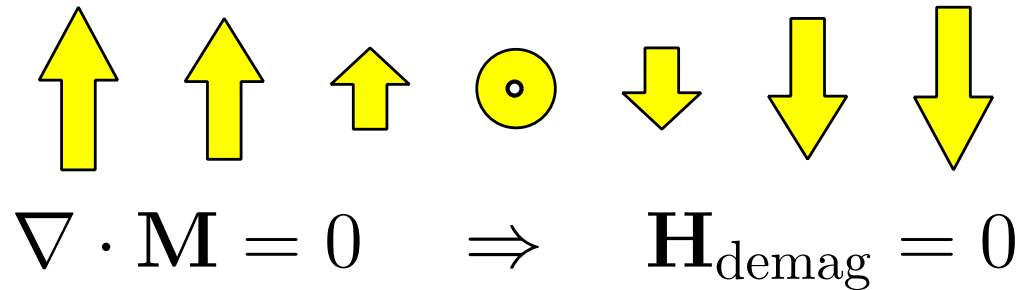
Subsample convergence:



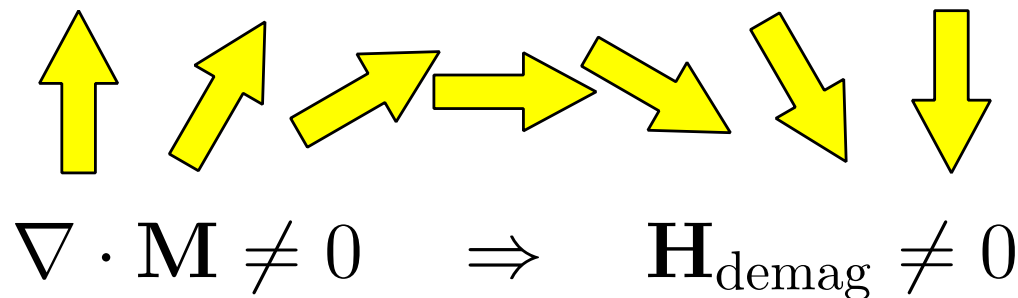


# Wall types

## Bloch wall

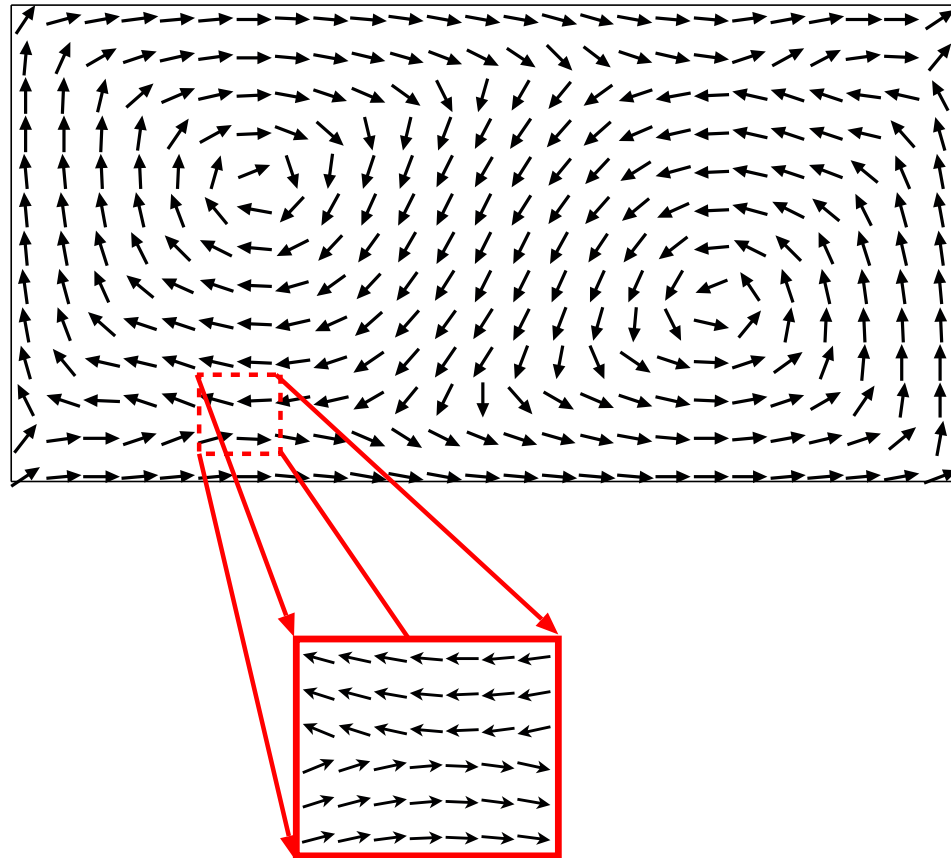


## Néel wall



# Néel-wall collapse

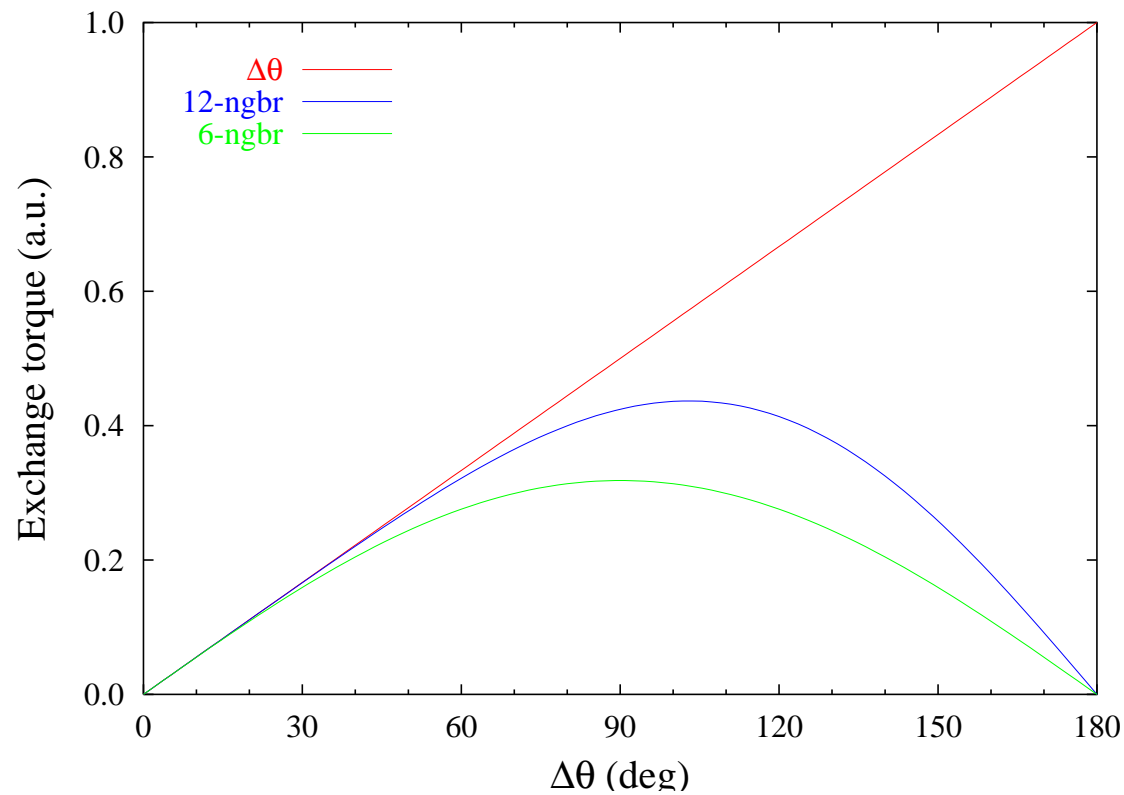
6-pt exchange,  $\mu_0 H = 5 \text{ mT}$ ,  $h = 20 \text{ nm}$



# Magnetization spiral

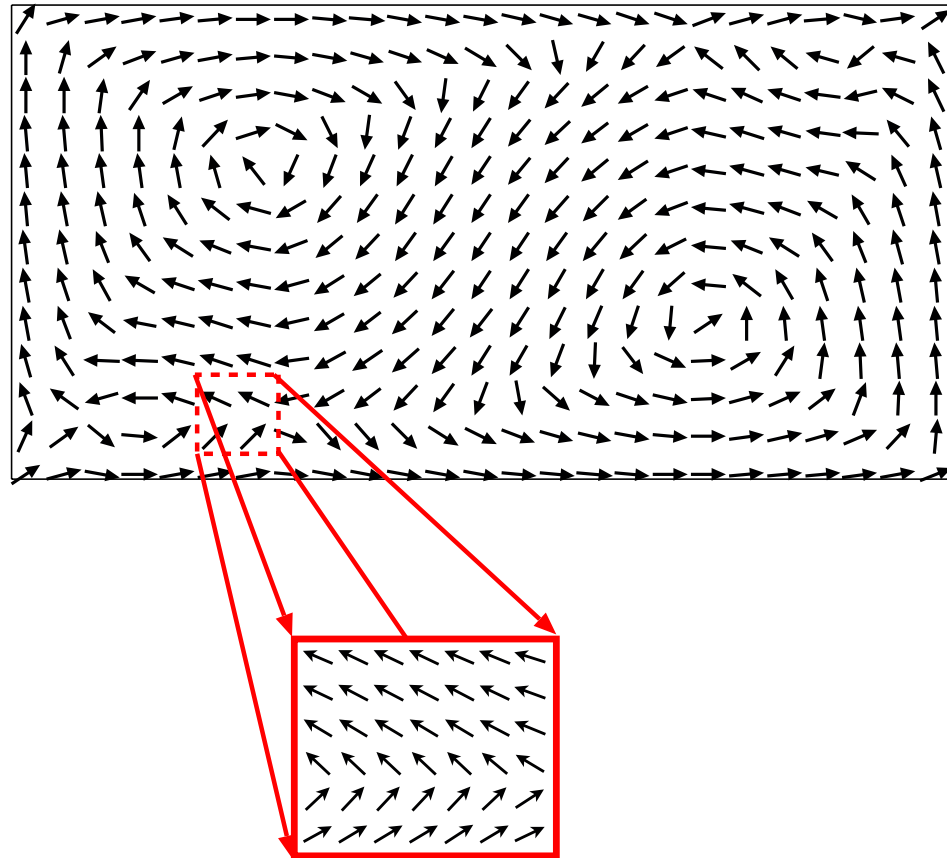
$$\mathbf{m} = (\cos \omega x, \sin \omega x)$$

Exchange torque vs.  $\omega$



# *Néel-wall non-collapse*

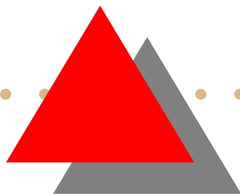
12-pt exchange,  $\mu_0 H = 6 \text{ mT}$ ,  $h = 20 \text{ nm}$





# Conclusions

- 6-ngbr and 26-ngbr are 2nd order.
- 12-ngbr is 4th order.
- Proper boundary conditions must be applied.
- Eigenvalues determine iterative behavior.
- 26-ngbr has less pinning for large cells, 12-ngbr dominates for  $h < l_{\text{ex}}$ .
- 12-ngbr helps against Néel wall collapse.





# References

- $\mu$ MAG:

<http://www.ctcms.nist.gov/~rdm/mumag.org.html>

- OOMMF:

<http://math.nist.gov/oommf>

- In press:

M.J. Donahue and D.G. Porter, “Exchange Energy Formulations for 3D Micromagnetics,”  
*Physica B*