

Quantum Algorithms for Quantum Field Theories

Stephen Jordan

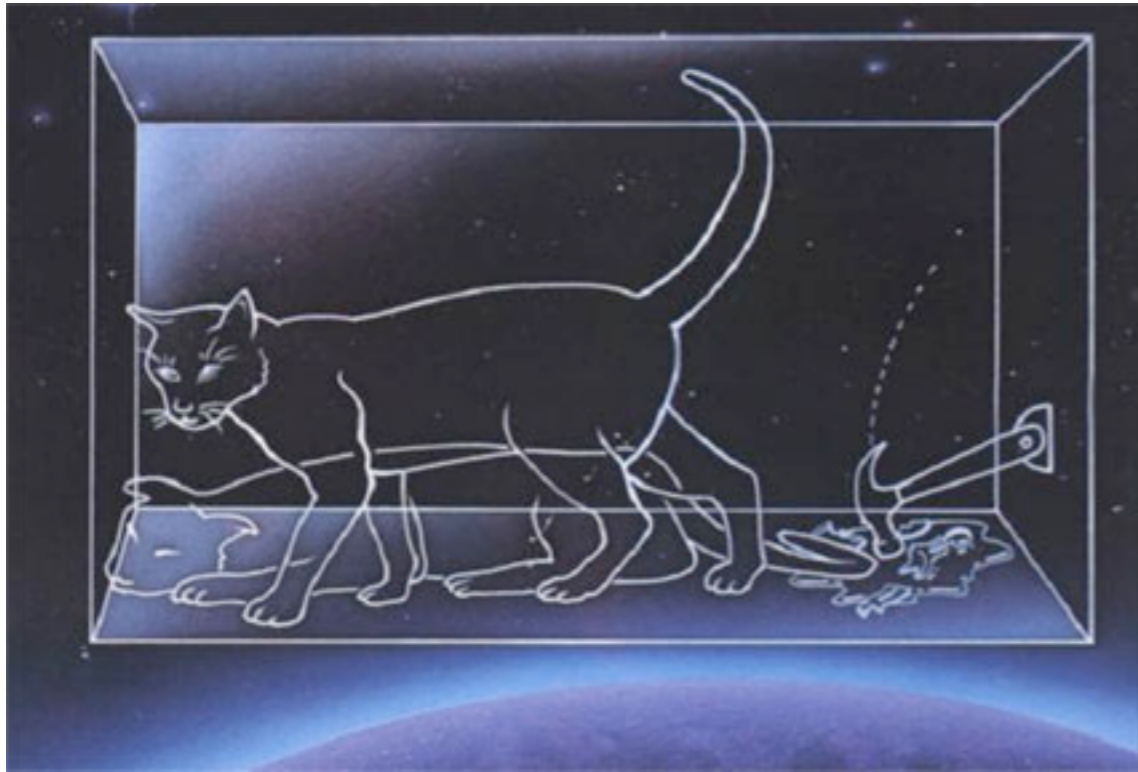
Joint work with

Keith Lee

John Preskill

[arXiv:1111.3633 and 1112.4833]

Quantum Mechanics



Each state of the system is a basis vector.

$|\text{dead}\rangle$

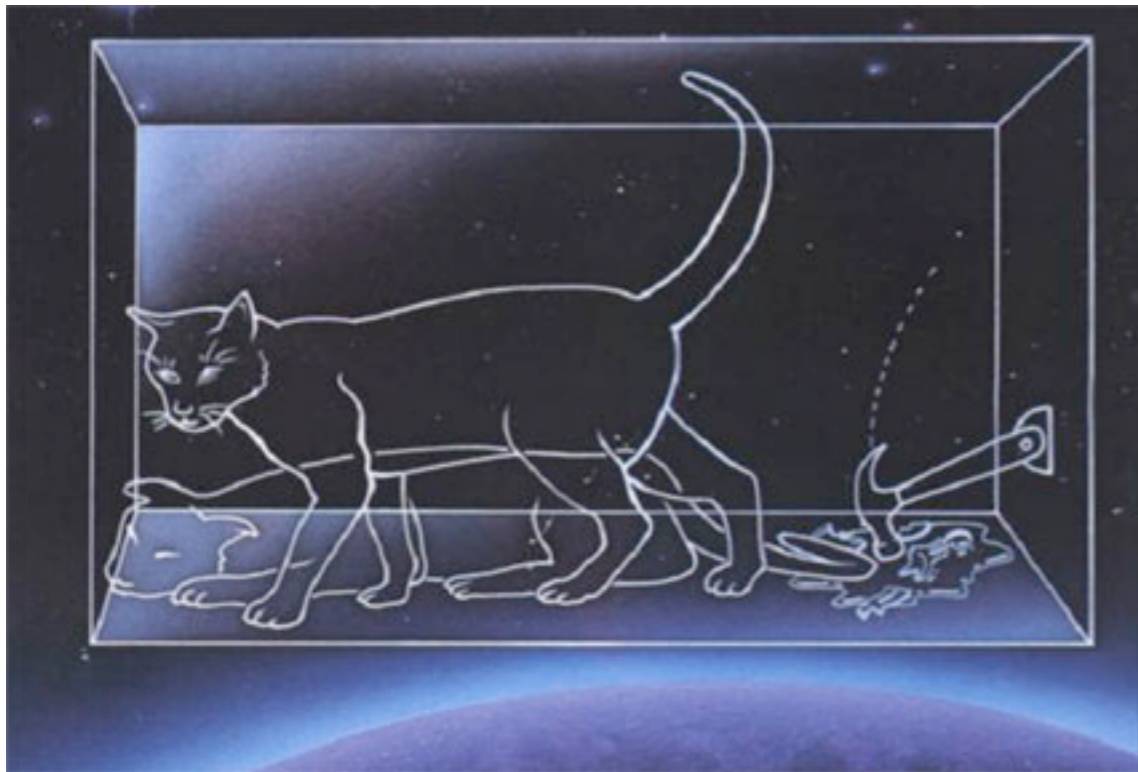
$|\text{alive}\rangle$

A general state is a linear combination of this basis:

$$\alpha|\text{dead}\rangle + \beta|\text{alive}\rangle$$

$$\alpha, \beta \in \mathbb{C}$$

Quantum Mechanics



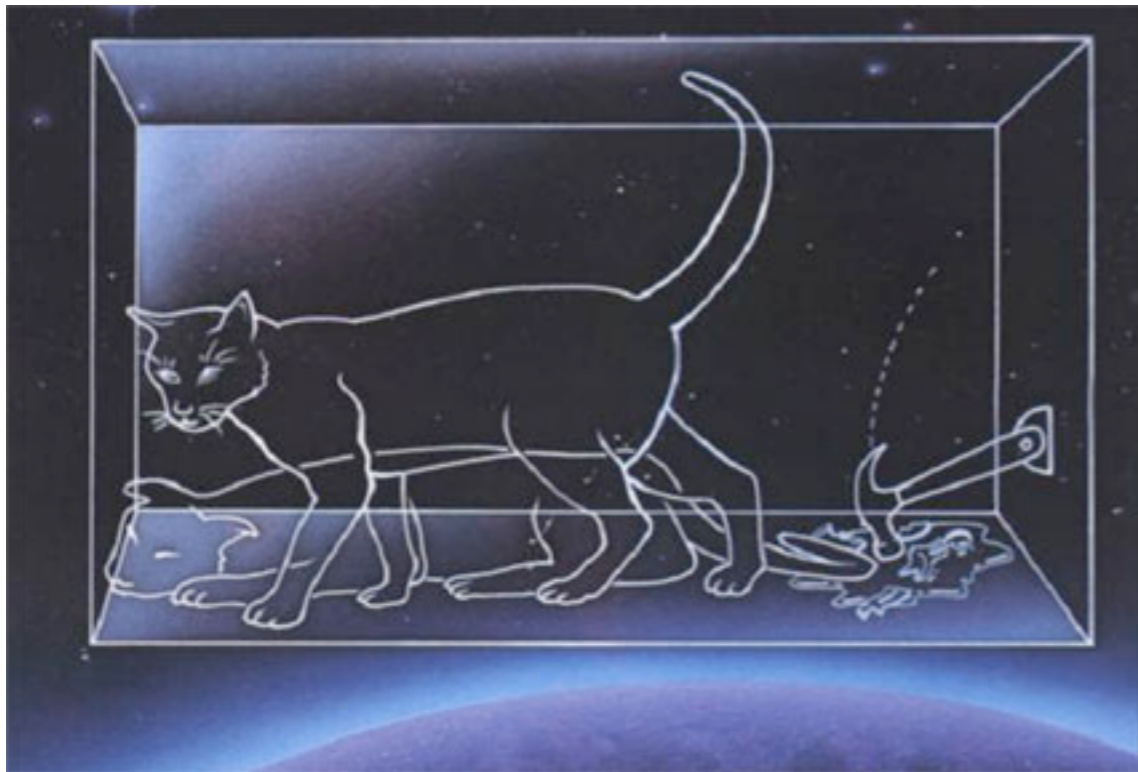
$$\alpha|\text{dead}\rangle + \beta|\text{alive}\rangle$$

If we look inside the box we see:

A dead cat with probability $|\alpha|^2$

A living cat with probability $|\beta|^2$

The Classical World



In most macroscopic systems, noise from the environment randomizes the phases.

The linear combination of states then acts like an ordinary probability distribution.

$$(p_{\text{dead}}, p_{\text{alive}}) \in \mathbb{R}^2$$

Qubits

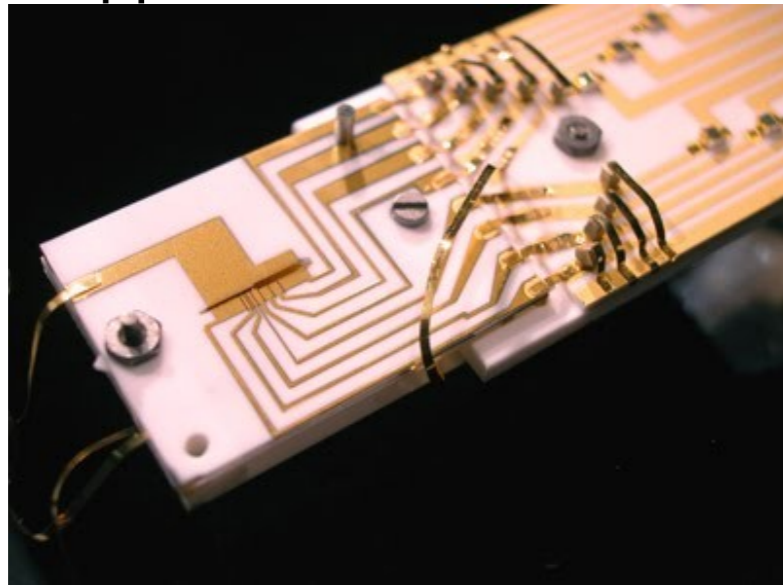
To exhibit quantum-mechanical effects we want a system that is simple and well isolated from its environment.

One qubit: $\alpha|0\rangle + \beta|1\rangle$

n qubits: $\sum_{x \in \{0,1\}^n} \alpha(x)|x\rangle$

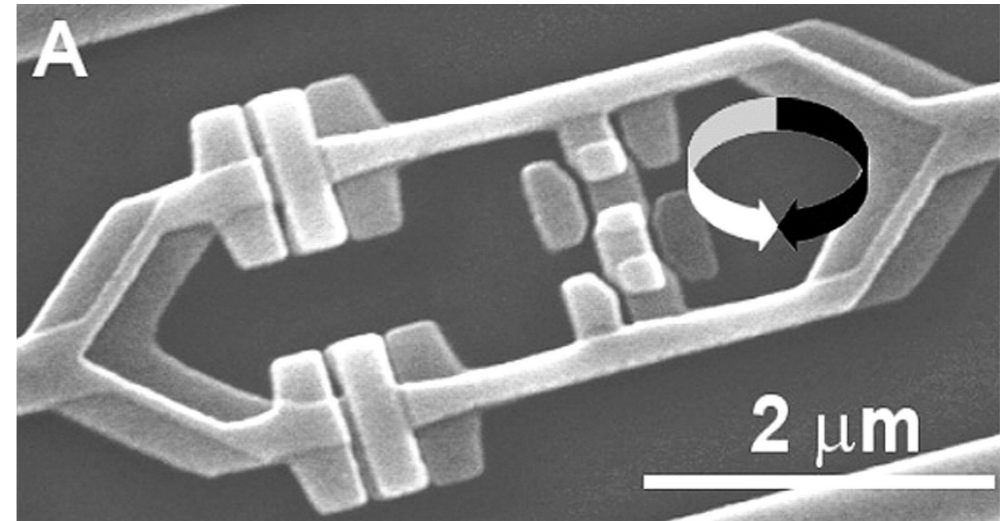
Qubits

Trapped Ions



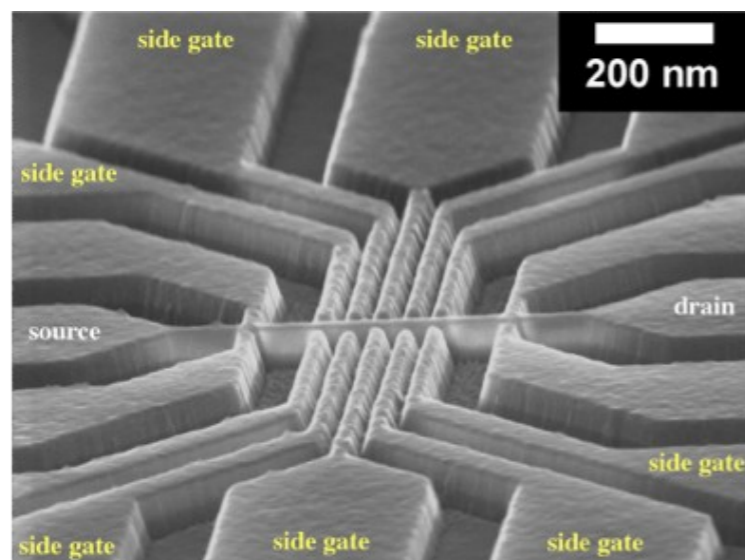
[Wineland group, NIST]

Superconducting Circuits



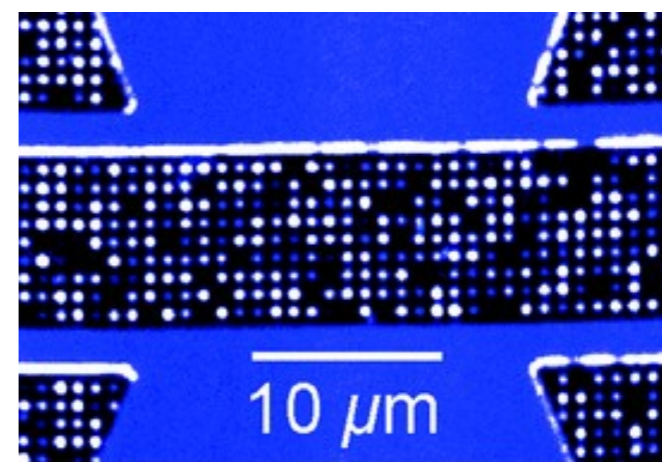
[Mooij group, TU Delft]

Quantum Dots



[Paul group, U. Glasgow]

NV Centers in Diamond

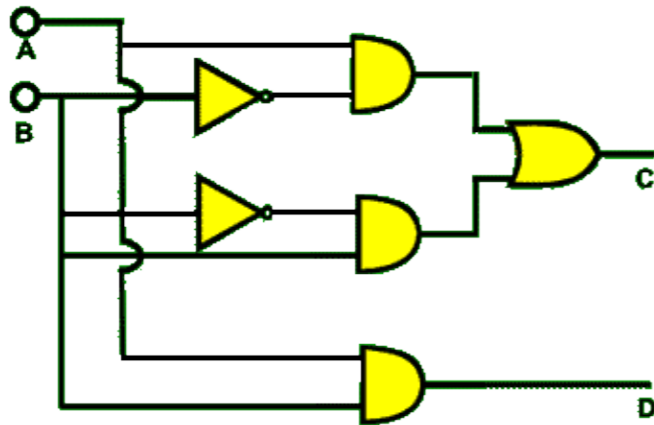


[Awschalom group, UCSB]

Quantum Circuits

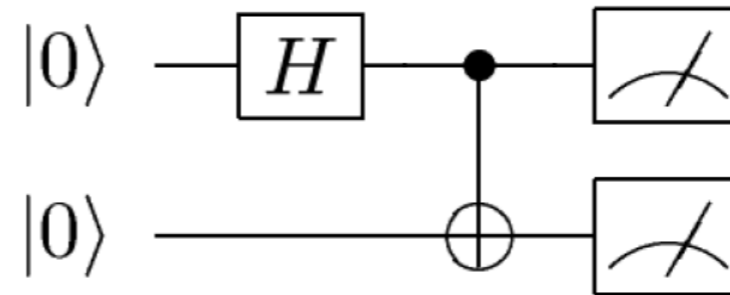
Classical

0101101

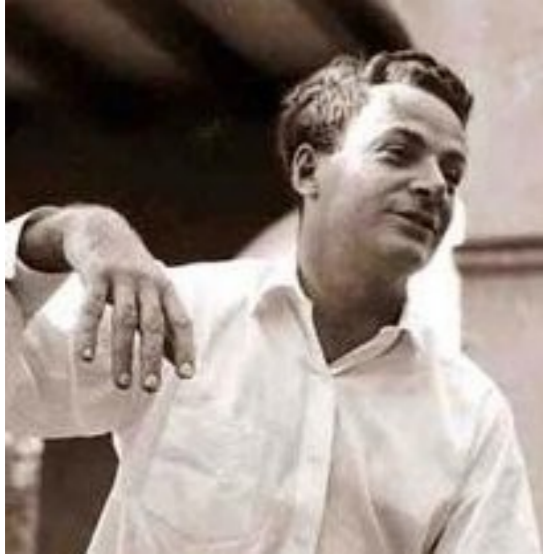


Quantum

$$|\psi\rangle = \sum_{x \in \{0,1\}^n} \alpha(x) |x\rangle$$



$$\text{---} \boxed{H} \text{---} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$



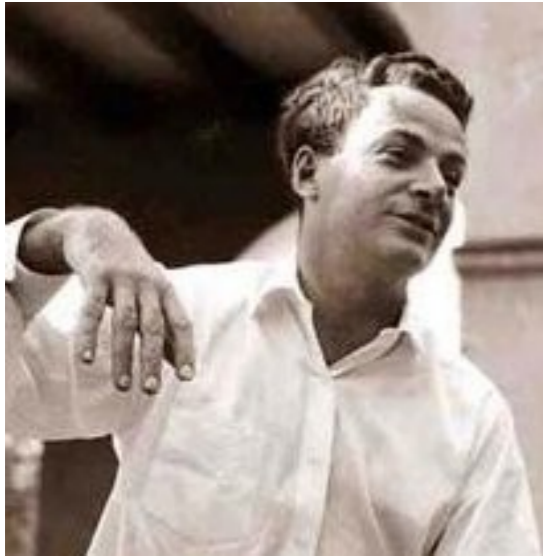
The full description of quantum mechanics for a large system with R particles has too many variables. It cannot be simulated with a normal computer with a number of elements proportional to R .

-Richard Feynman, 1982



An n -bit integer can be factored on a quantum computer in $O(n^2)$ time.

-Peter Shor, 1994



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Are there any systems that remain hard to simulate even with quantum computers?

Quantum Simulation

Condensed-matter lattice models:

[Lloyd, 1996]

[Abrams, Lloyd, 1997]

[Berry, Childs, 2012]

Many-particle Schrödinger and Dirac Equations:

[Meyer, 1996]

[Zalka, 1998]

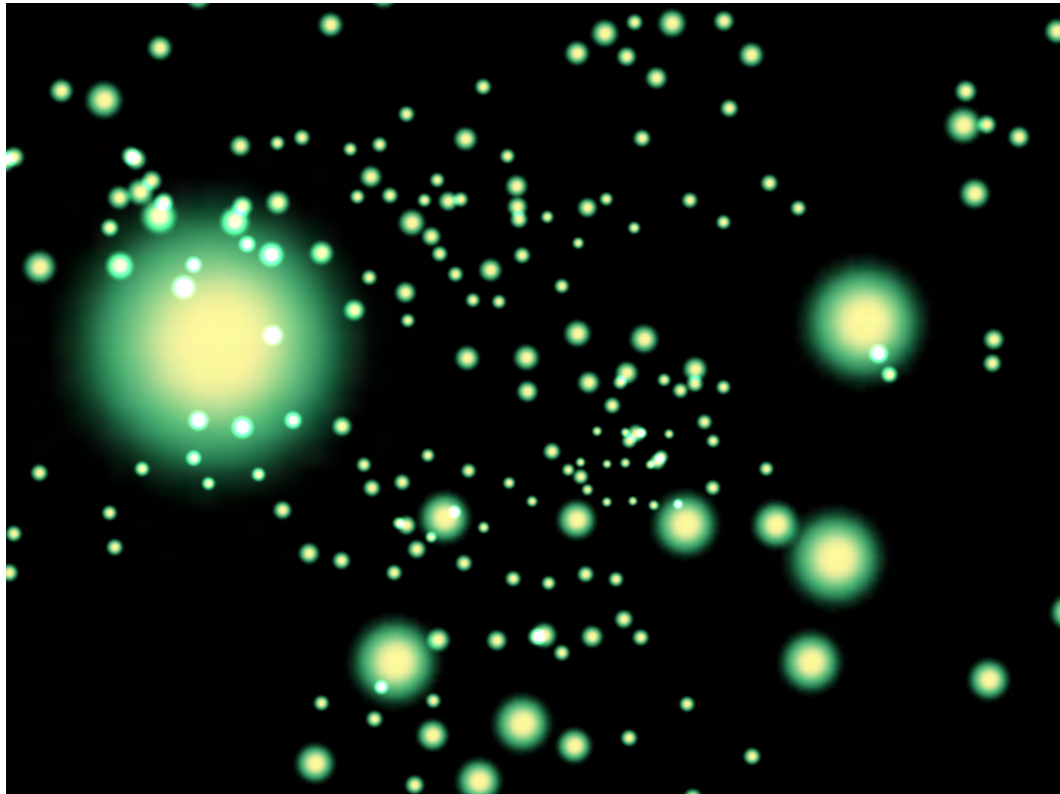
[Taylor, Boghosian, 1998]

[Kassal, S.J., Love, Mohseni, Aspuru-Guzik, 2008]

Quantum Field Theory

- Much is known about using quantum computers to simulate quantum systems.
- Why might QFT be different?
 - Field has infinitely many degrees of freedom
 - Relativistic
 - Particle number not conserved
 - Formalism looks different

Quantum Particles



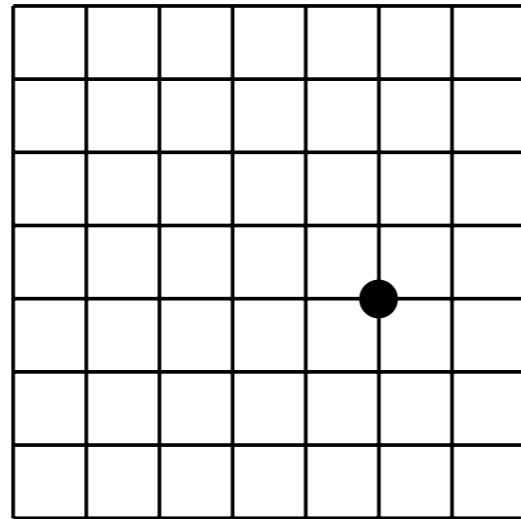
A classical particle is described by its location coordinates.

$$\vec{r} = (x, y, z)$$

The state of a quantum particle is linear combination of positions.

$$|\psi\rangle = \int d^3r \psi(r) |r\rangle$$

A configuration is a list of particle coordinates.

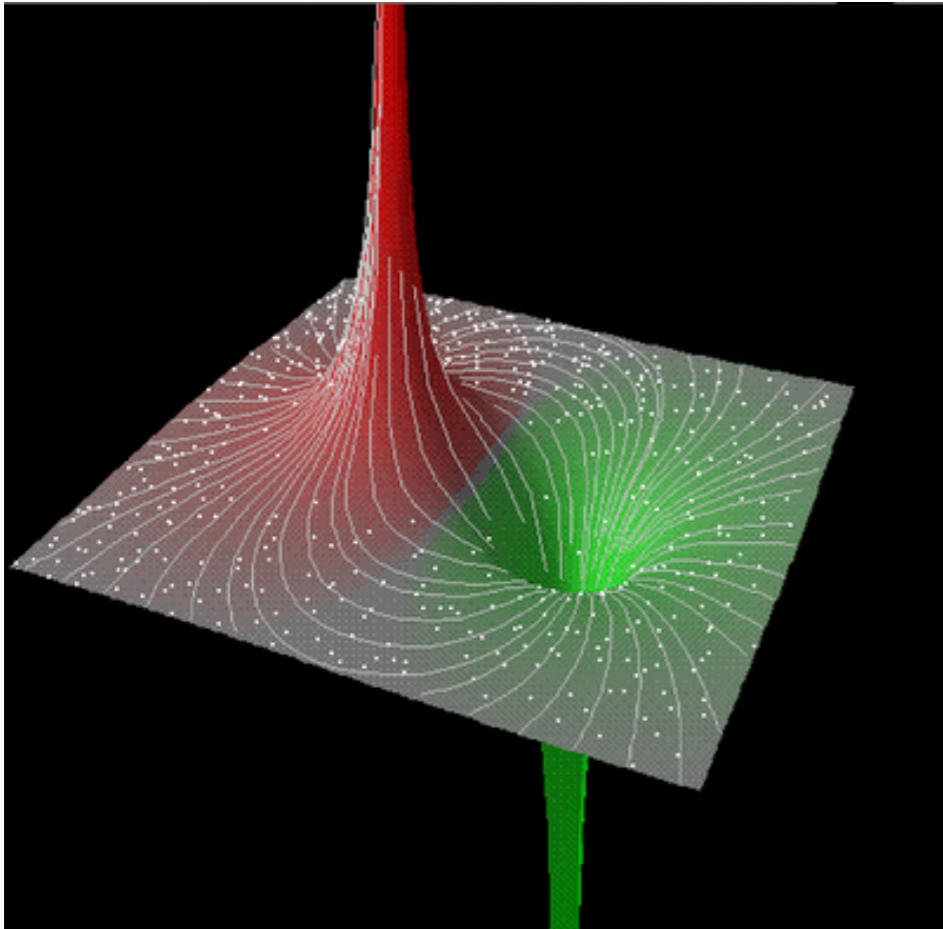


(5,3)

A quantum particle can be in a superposition of locations.

$$\frac{1}{\sqrt{2}} \left| \begin{array}{c} \text{7x7 grid with dot at (5,3)} \\ (5,3) \end{array} \right\rangle - \frac{i}{\sqrt{2}} \left| \begin{array}{c} \text{7x7 grid with dot at (2,2)} \\ (2,2) \end{array} \right\rangle$$

Quantum Fields



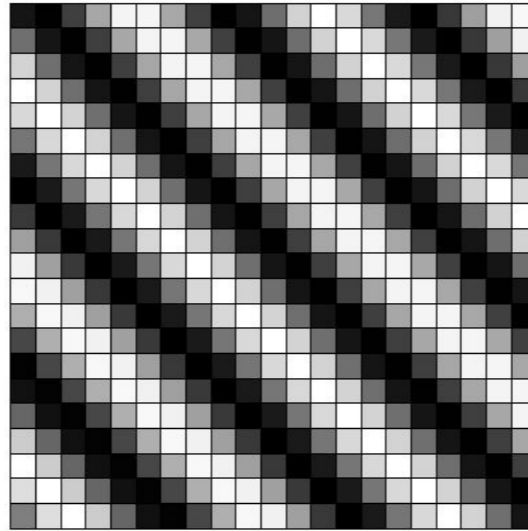
A classical field is described by its value at every point in space.

$$E(r) = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2}$$

A quantum field is a linear combination of classical field configurations.

$$|\Psi\rangle = \int \mathcal{D}[E] \Psi[E] |E\rangle$$

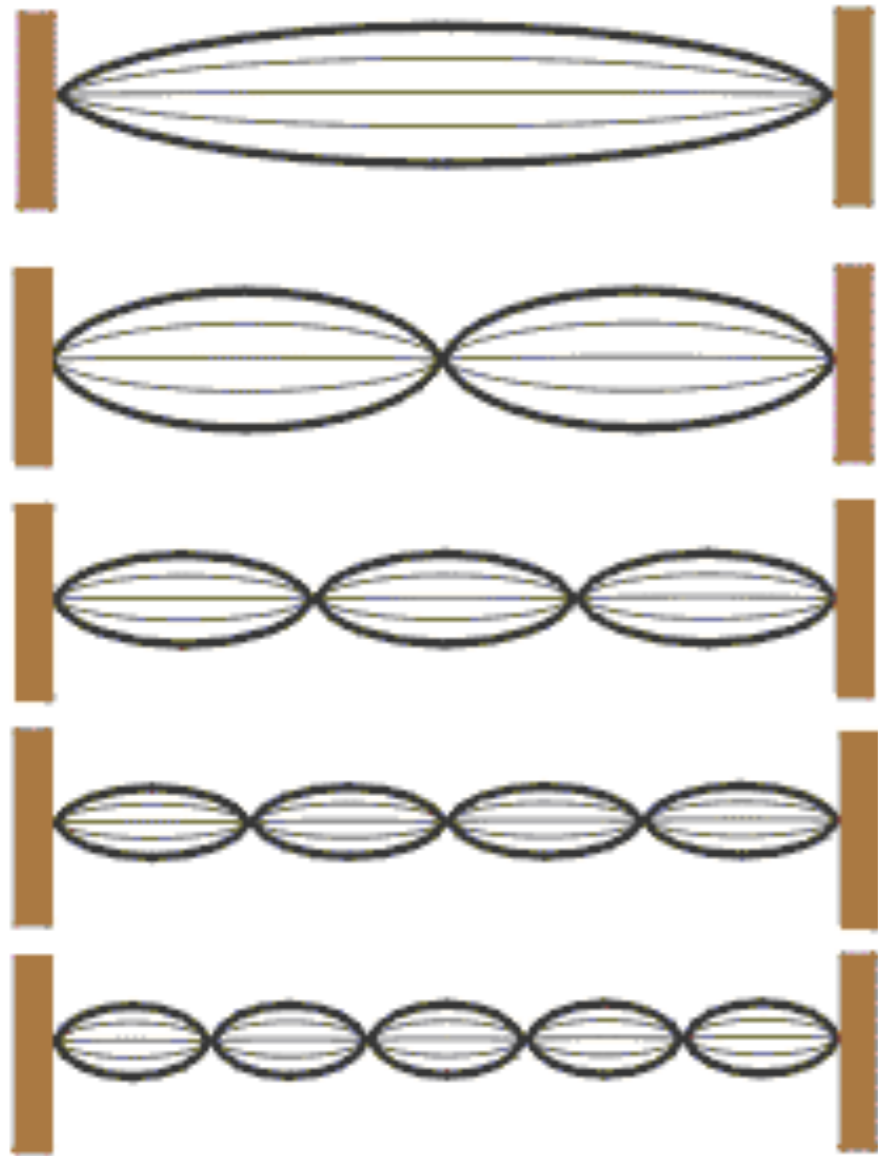
A configuration of the field is a list of field values, one for each lattice site.



A quantum field can be in a superposition of different field configurations.

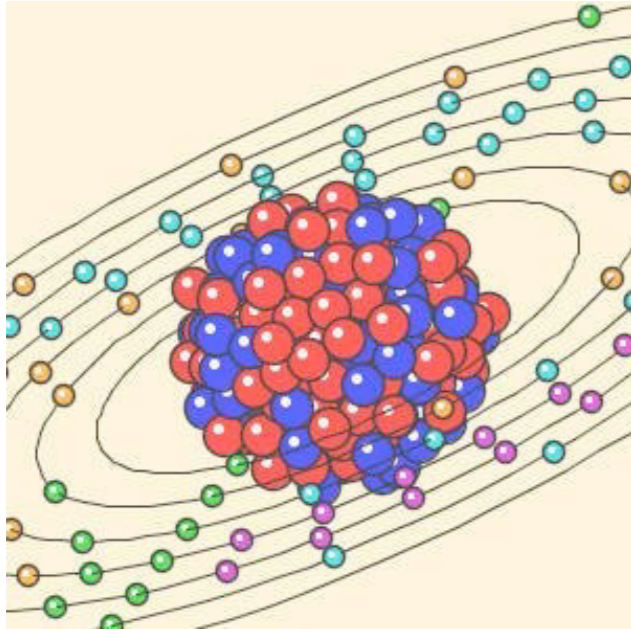
$$\frac{1}{\sqrt{2}} \left| \begin{array}{c} \text{[Smooth Gradient Grid]} \end{array} \right\rangle - \frac{i}{\sqrt{2}} \left| \begin{array}{c} \text{[Diagonal Pattern Grid]} \end{array} \right\rangle$$

Particles Emerge from Fields

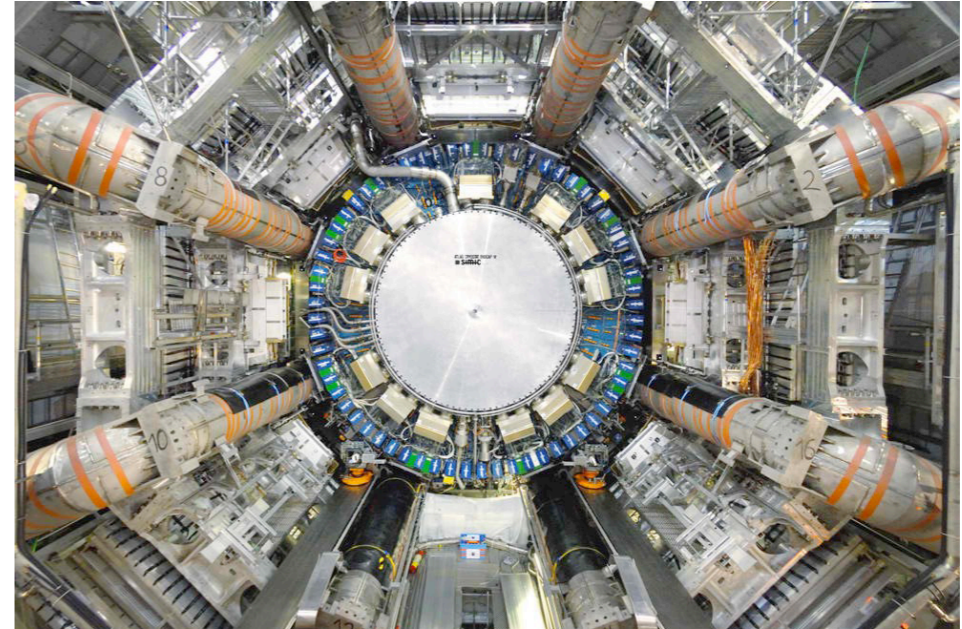


Particles of different energy are different resonant excitations of the field.

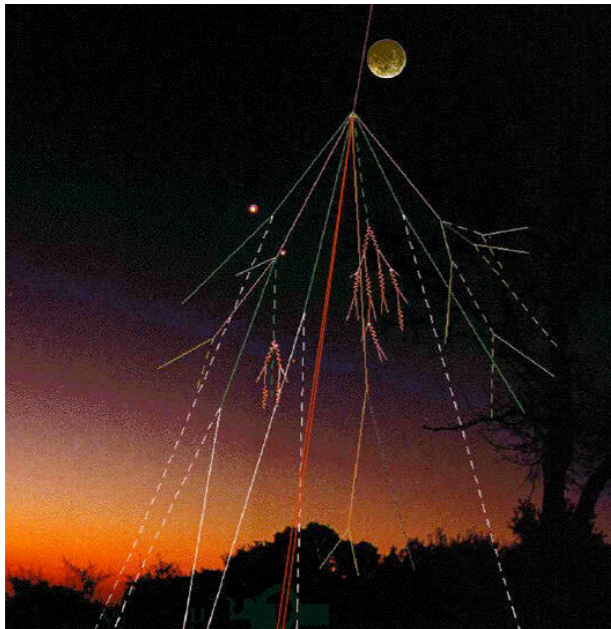
When do we need QFT?



Nuclear Physics



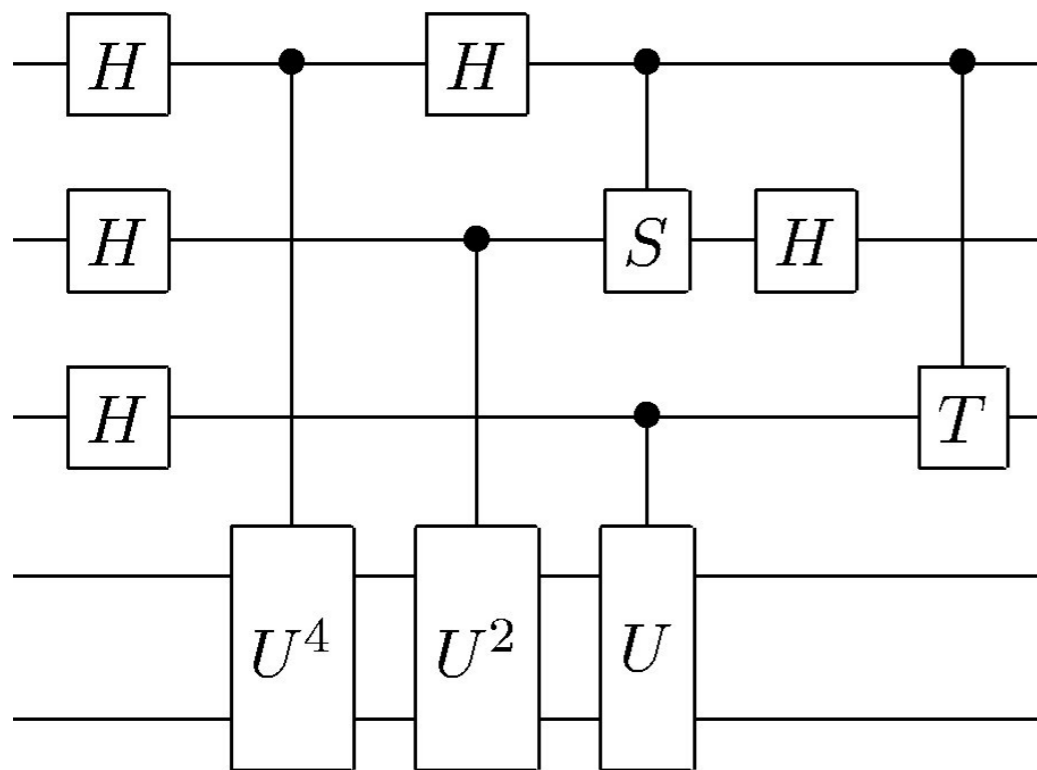
Accelerator Experiments



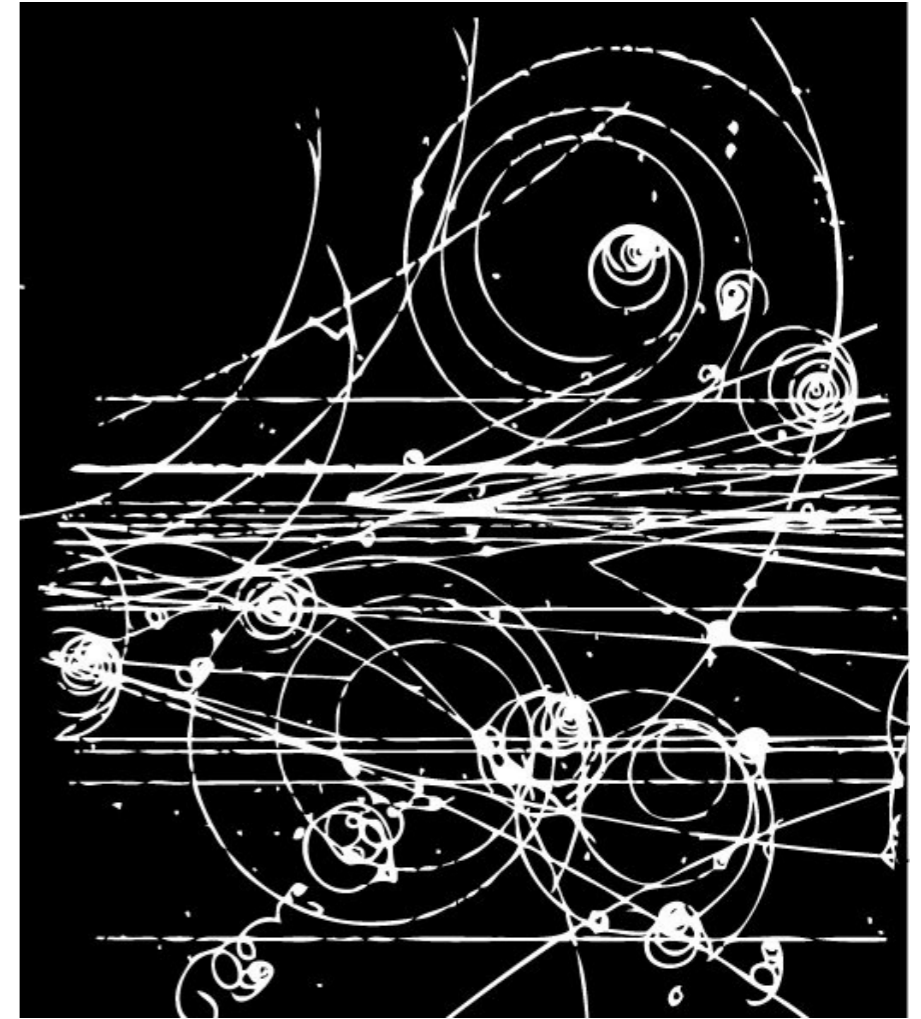
Cosmic Rays

→Whenever quantum mechanical and relativistic effects are both significant.

What is the computational power of our universe?

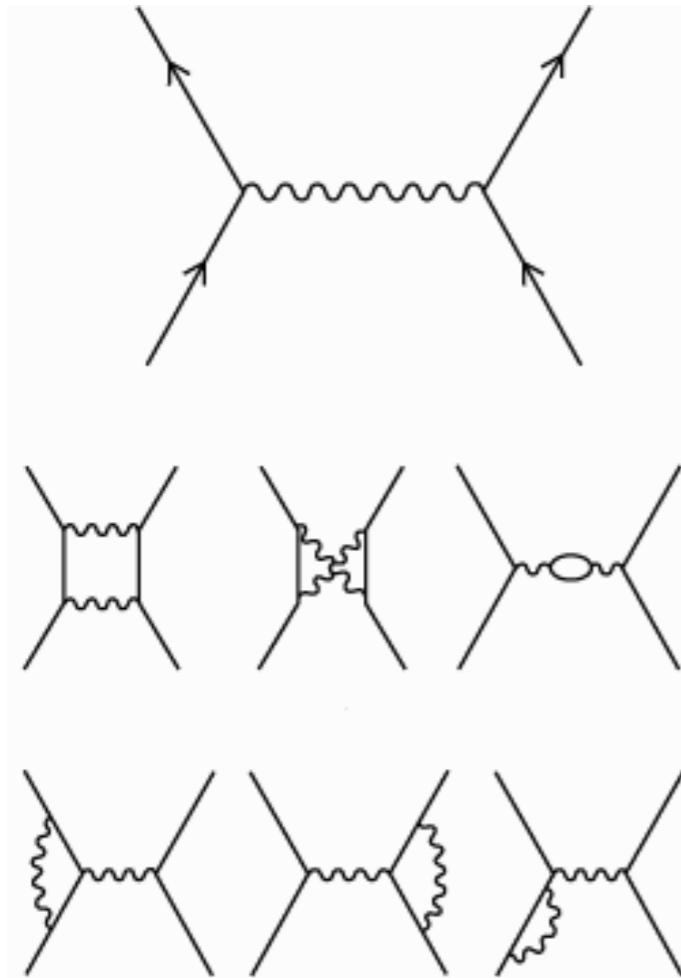


→
simulate
←



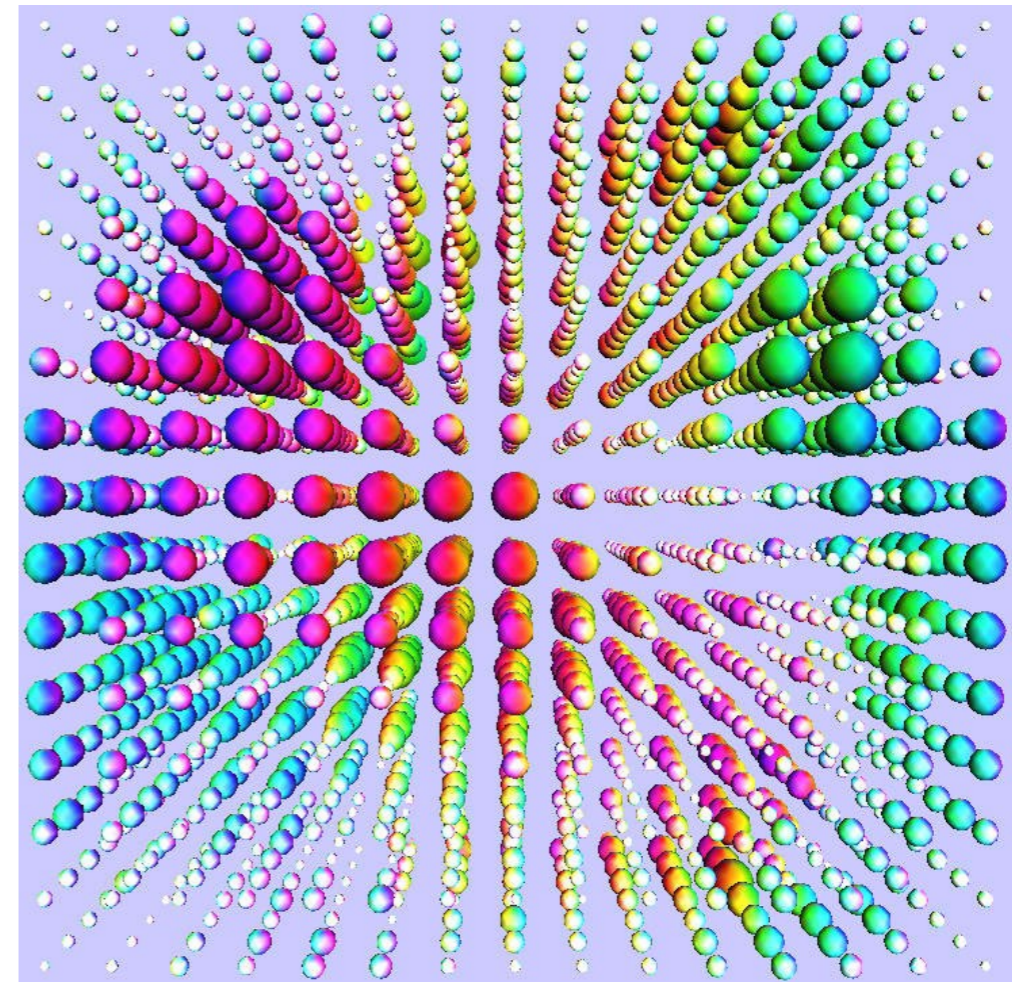
Classical Algorithms

Feynman diagrams



Break down at strong coupling or high precision

Lattice methods

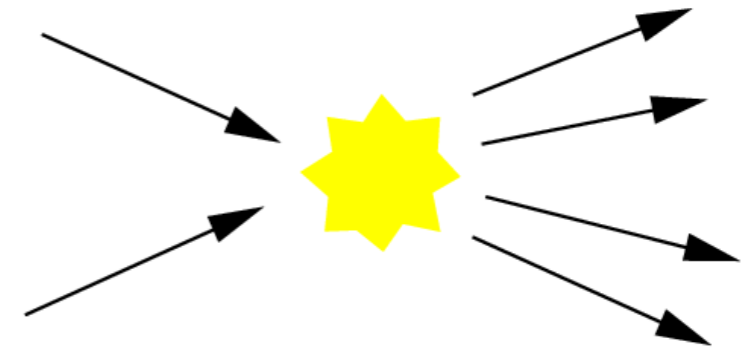


Cannot calculate scattering amplitudes

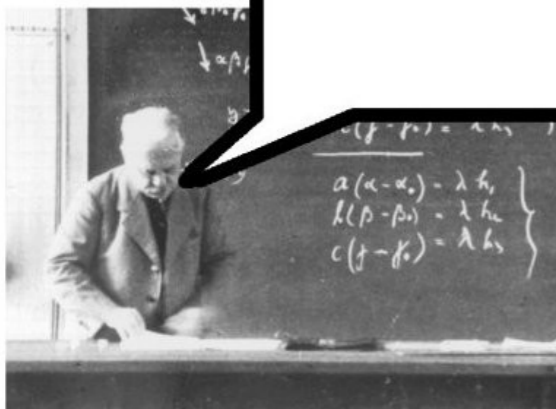
A QFT Computational Problem

Input: a list of momenta of incoming particles

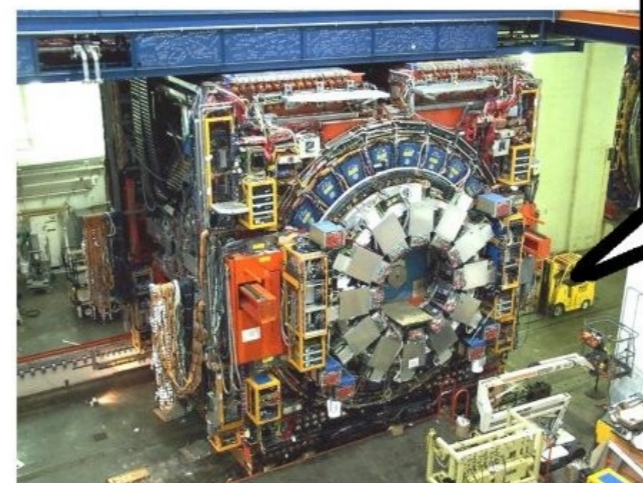
Output: a list of momenta of outgoing particles



S-matrix



Particle
accelerator



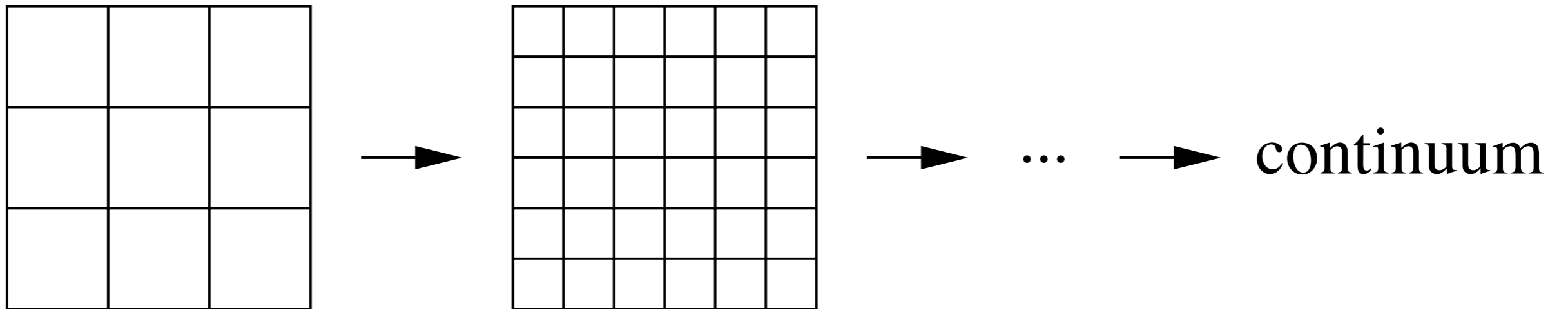
I will present a polynomial-time quantum algorithm to compute scattering probabilities in ϕ^4 -theory with nonzero mass

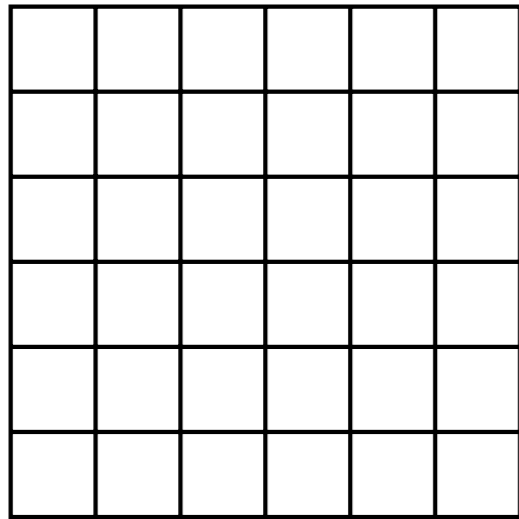
ϕ^4 -theory is a simple model that illustrates some of the main difficulties in simulating a QFT:

- Discretizing spacetime
- Preparing initial states
- Measuring observables

Lattice cutoff

Continuum QFT = limit of a sequence of theories on successively finer lattices



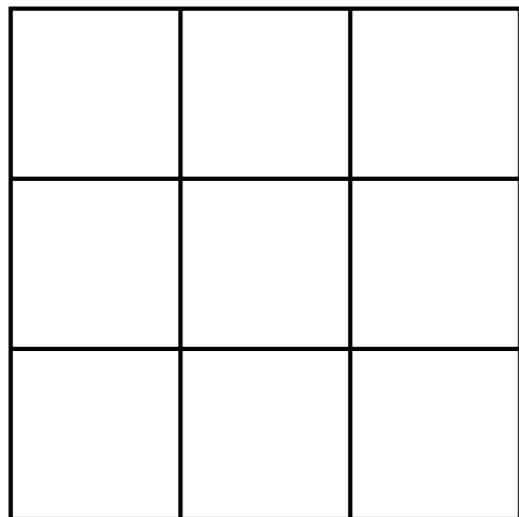


Mass: m

Interaction strength: λ



Coarse grain

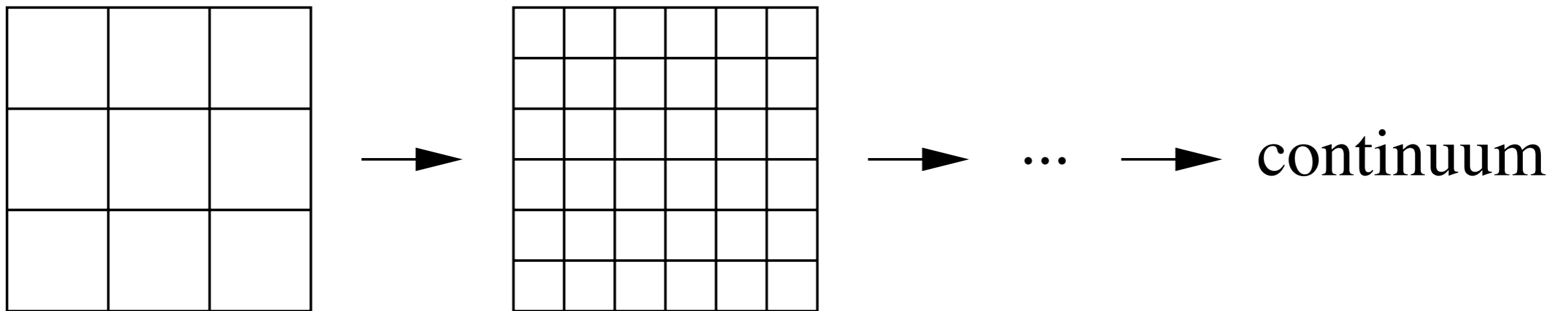


Mass: m'

Interaction strength: λ'

Lattice cutoff

Continuum QFT = limit of a sequence of theories on successively finer lattices



m and λ are functions of lattice spacing!

Discretization Errors

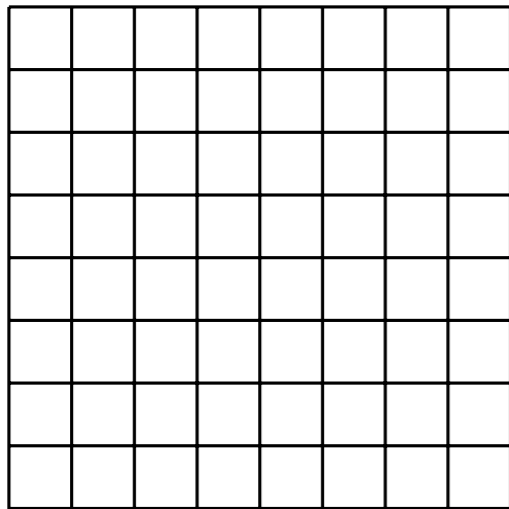
- Renormalization of m and λ make discretization tricky to analyze
- In ϕ^4 -theory, in $d=1,2,3$, discretization errors scale as a^2

$$\begin{aligned}
 \text{---} \bigcirc \text{---} &= \frac{(-i\lambda_0)^2}{6} \iint \frac{d^D k}{(2\pi)^D} \frac{d^D q}{(2\pi)^D} \frac{i}{(k^0)^2 - \sum_i \frac{4}{a^2} \sin^2 \left(\frac{ak^i}{2} \right) - m^2} \frac{i}{(q^0)^2 - \sum_i \frac{4}{a^2} \sin^2 \left(\frac{aq^i}{2} \right) - m^2} \\
 &\quad \times \frac{i}{(p^0 + k^0 + q^0)^2 - \sum_i \frac{4}{a^2} \sin^2 \left(\frac{a(p^i + k^i + q^i)}{2} \right) - m^2} \quad (207)
 \end{aligned}$$

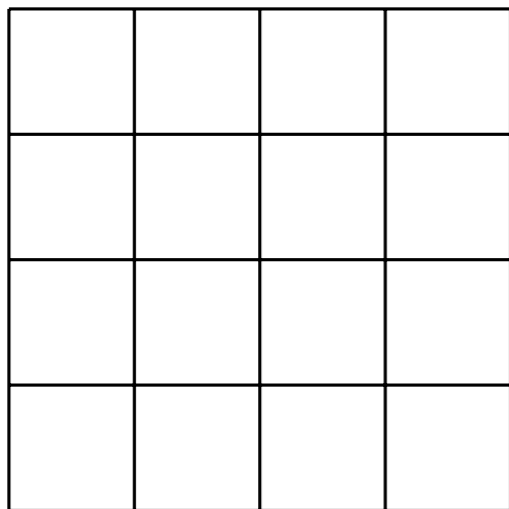
$$= \frac{i\lambda_0^2}{3} \int_0^1 \int_0^1 \int_0^1 dx dy dz \delta(x + y + z - 1) \iint \frac{d^D k}{(2\pi)^D} \frac{d^D q}{(2\pi)^D} \frac{1}{D^3}, \quad (208)$$

...its complicated

Condensed Matter



↓ RG



There is a fundamental lattice spacing.

But:

We may save qubits by simulating a coarse-grained theory.

After imposing a spatial lattice we have a many-body quantum system with a local Hamiltonian

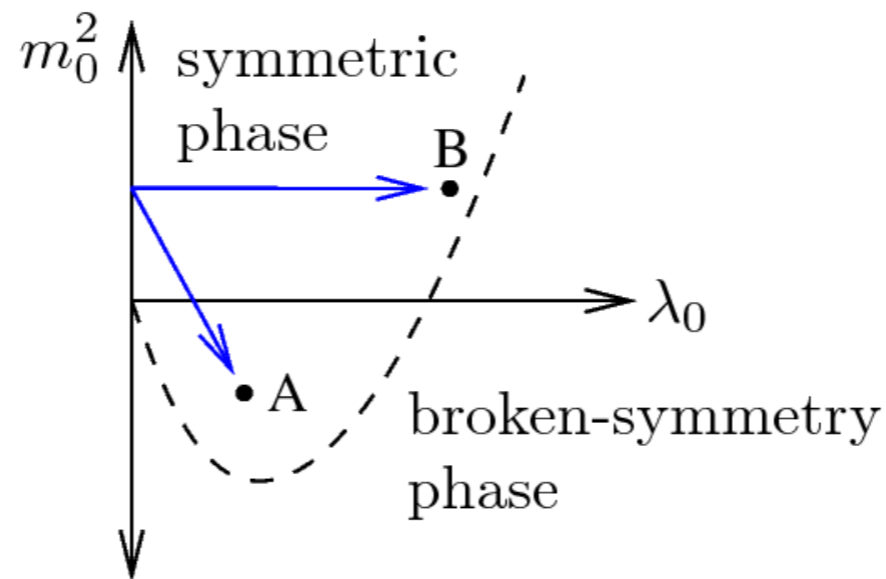
Simulating the time evolution in polynomial time is a **solved problem**

Standard methods scale as N^2 . We can do N .

- **Convergence as $a \rightarrow 0$**
- Preparing wavepackets
- Measuring particle momenta

Strong Coupling

ϕ^4 -theory in 1+1 and 2+1 dimensions has a quantum phase transition in which the $\phi \rightarrow -\phi$ symmetry is spontaneously broken



Near the phase transition perturbation theory fails and the gap vanishes.

$$m_{\text{phys}} \sim (\lambda_c - \lambda_0)^\nu \quad \nu = \begin{cases} 1 & d = 1 \\ 0.63 \dots & d = 2 \end{cases}$$

Complexity

Weak Coupling:

| | |
|---------|------------------------|
| $d = 1$ | $(1/\epsilon)^{1.5}$ |
| $d = 2$ | $(1/\epsilon)^{2.376}$ |
| $d = 3$ | $(1/\epsilon)^{5.5}$ |

Strong Coupling:

| | $\lambda_c - \lambda_0$ | p | n_{out} |
|---------|---|-------|--------------------------|
| $d = 1$ | $\left(\frac{1}{\lambda_c - \lambda_0}\right)^8$ | p^4 | n_{out}^5 |
| $d = 2$ | $\left(\frac{1}{\lambda_c - \lambda_0}\right)^{5.04}$ | p^6 | $n_{\text{out}}^{7.128}$ |

Eventual goal:

Simulate the standard model in BQP

Solved problems:

ϕ^4 -theory [arXiv:1111.3633 and 1112.4833]

Gross-Neveu [S.J., Lee, Preskill, *in preparation*]

Open problems:

Gauge symmetries, massless particles

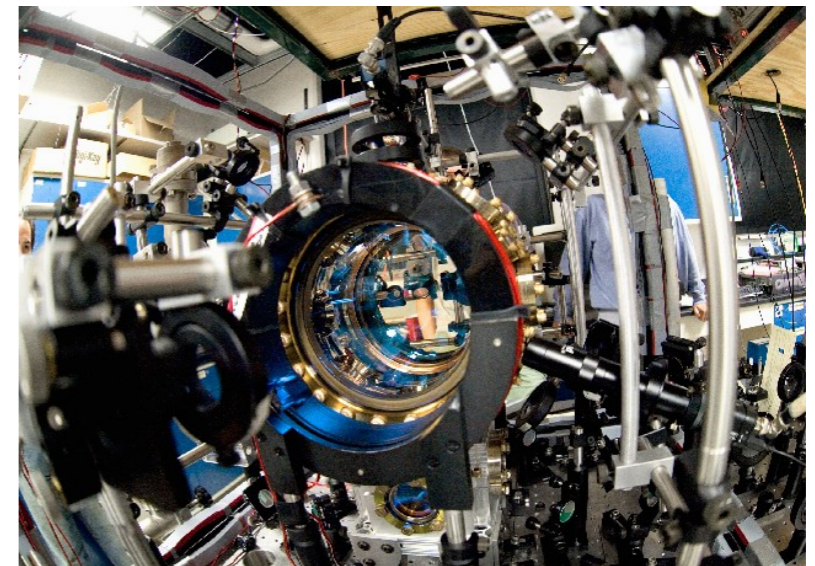
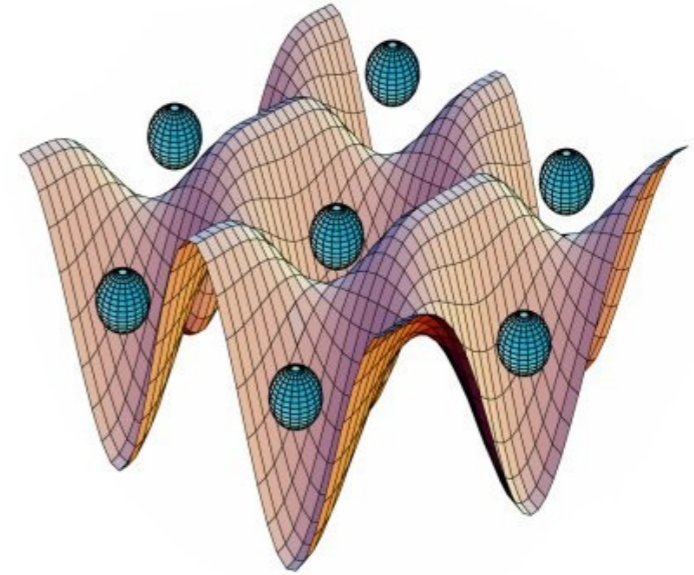
Spontaneous symmetry breaking

Bound states, confinement

Chiral Fermions

Analog Simulation

- No gates: just implement a Hamiltonian and let it time-evolve
- Current experiments do this!

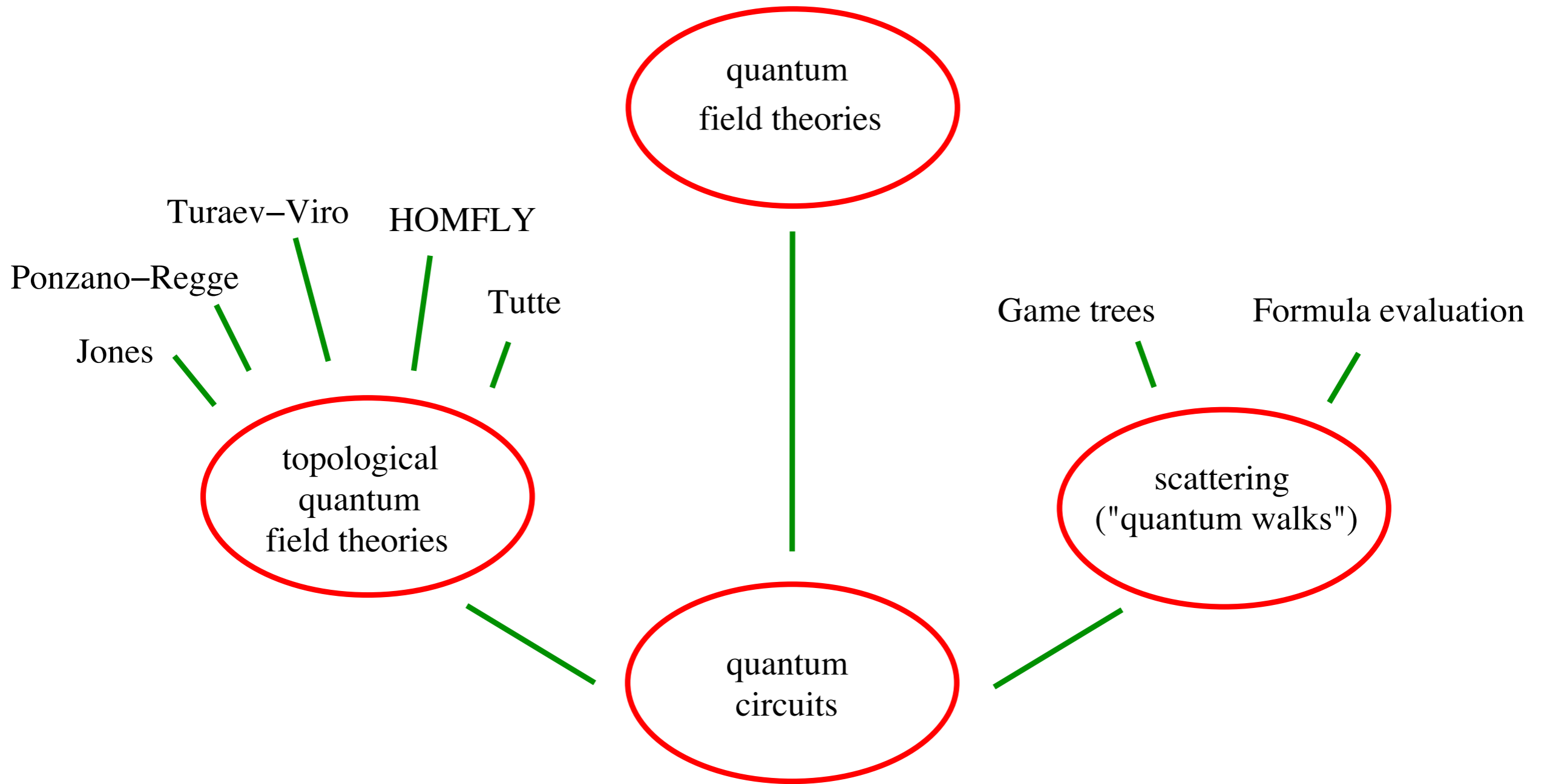


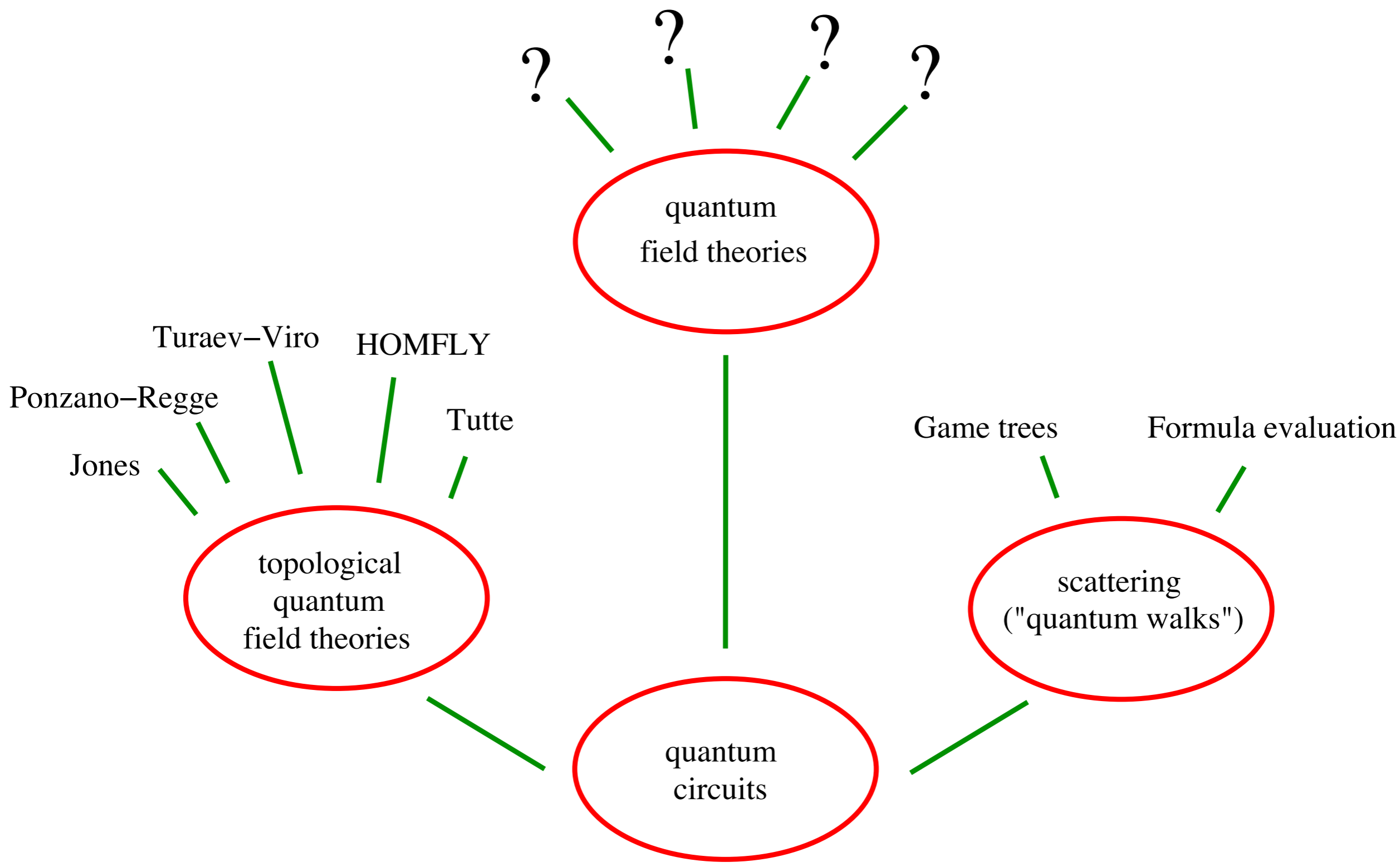
Analog Simulation

- Experiments so far have concentrated on mapping out phase diagrams
- We are developing a proposal to simulate ϕ^4 scattering processes using Rydberg atoms trapped in optical lattices

[Gorshkov, S.J., Preskill, Lee, *In Preparation*]

Broader Context







What I'm trying to do is get you people who think about computer simulation to see if you can't invent a different point of view than the physicists have.

-Richard Feynman, 1981



In thinking and trying out ideas about “what is a field theory” I found it very helpful to demand that a correctly formulated field theory should be soluble by computer.. It was clear, in the '60s, that no such computing power was available in practice.

-Kenneth Wilson, 1982

Conclusion

Quantum computers can simulate scattering in ϕ^4 -theory.

There are many exciting prospects for quantum computation and quantum field theory to contribute to each other's progress.

I thank my collaborators:



Thank you for your attention.